Experimentally-Based Initialization of Direct Numerical Simulations of Miscible, Rayleigh-Taylor Instability-Induced Mixing

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The objective of this research is to couple experimentallymeasured initial conditions with high-resolution DNS

- Experimental Objectives
 - Measure initial density and velocity conditions of a small Atwood number, Rayleigh-Taylor driven mixing layer[†]
 - Quantify perturbations in both streamwise and cross-stream directions
 - Measure time-evolution of density statistics, spectra, and molecular mixing parameters
- Numerical Objectives
 - Develop and validate method of initializing DNS (direct numerical simulation) with experimentally-measured data
 - Perform 2D and 3D DNS using measured initial conditions for comparison with experiment[‡] and compare 2D and 3D dynamics
 - Investigate effect of initial conditions on transition from linear to nonlinear mixing layer evolution
- [†] P. Ramaprabhu & M. J. Andrews, "Experimental investigation of Rayleigh–Taylor mixing at small Atwood numbers," J. Fluid Mech. 502, 233 (2004)
- [‡] P. Ramaprabhu & M. J. Andrews, "On the initialization of Rayleigh-Taylor simulations," Phys. Fluids 16, L59 (2004)

Taylor's hypothesis is used to relate the temporal evolution of statistics to the spatial evolution

 ρ_1

- $\Lambda T \approx 5^{\circ}$
- $A = 7.5 \times 10^{-4}$
- $U_m \approx 4.2$ cm/s mean velocity
- Pr = 7
- Downstream distance x from splitter plate related to time by

$$t = \frac{x}{U_m}$$

Dimensionless time

$$\tau = t \sqrt{\frac{gA}{H}} = \frac{x}{U_m} \sqrt{\frac{gA}{H}}$$



Density measurements at several downstream locations quantify the evolution of mixing statistics and PLIF diagnostics measured cross-stream perturbation



- Thermocouple measurements
 - Converted to density measurements through H₂0 equation of state
 - Smaller weld bead in new experiments decreased probe <u>volume</u> by 90%
- Planar Laser-Induced Fluorescence (PLIF) used to measure cross-stream (y-direction) perturbations off splitter plate
 - First experimental measurement of y-direction perturbation in Rayleigh-Taylor instability





The initial multi-mode interfacial perturbation $\zeta(x,y)$ is modeled by two Fourier series

Initial density assumed to have an error function profile

$$\rho(x,y,z,t=0) = \frac{\rho_1 + \rho_2}{2} + \frac{\rho_1 - \rho_2}{2} \operatorname{erf}\left[\frac{z + \zeta(x,y)}{\varepsilon}\right]$$

where $\varepsilon = \delta/2$ and δ is width of initial diffusion layer

$$\zeta(x, y) = \left[\sum_{m=0}^{N_x} a_m \cos\left(\frac{2\pi m}{\lambda_x}x\right) + b_m \sin\left(\frac{2\pi m}{\lambda_x}x\right)\right] + \left[\sum_{n=0}^{N_y} c_n \cos\left(\frac{2\pi n}{\lambda_y}y\right) + d_n \sin\left(\frac{2\pi n}{\lambda_y}y\right)\right]_{-1}$$

- Fourier coefficients taken directly from fluctuating density spectrum at x = 0.1 cm from splitter plate
- Initial diffusion velocity field required to satisfy continuity is

$$u_i = -\frac{D}{\rho} \frac{\partial \rho}{\partial x_i}$$

The initial density spectra in *x*- and *y*-directions include short- and long-wavelength perturbations

- Smallest perturbation present limited by current simulation resolutions
- Density spectra (thermocouple) and interfacial perturbation spectra (PLIF) related through interfacial thickness δ



Particle image velocimetry (PIV) was used to measure initial streamwise velocity perturbation

- Velocity measurements performed without buoyancy isolate momentum input generated by splitter plate
- Fluctuating vertical velocity (w') at x = 0.75 cm from edge of splitter plate characterize initial velocity conditions



The initial potential field is constructed from the measured initial vertical velocity spectrum



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• Potential field † $\phi(x, z) = a_0 \sum_{k=k_{\min}}^{k_{\max}} \frac{A(k)}{k} \sin(kx) \exp(-k|z|)$ created using spectral coefficients A(k)

from measured vertical velocity spectrum

Initial velocity defined as

$$u_i = \frac{\partial \phi}{\partial x_i} - \frac{D}{\rho} \frac{\partial \rho}{\partial x_i}$$

- Initial velocity conditions DNS have PLIF
 measured perturbation in *y*-direction and
 no ρ perturbation in *x*-direction
- Initial potential field for 2D and 3D DNS is 2D, as no velocity perturbation is measured in *y*-direction



[†] P. G. Drazin & W. H. Reid, *Hydrodynamic Stability*, Cambridge University Press (1982)

Simulations have been performed with both initial density and velocity conditions



- Four simulations have been performed:
 - Initial density conditions
 - 2D: 1024² (32 \times 32 cm)
 - 3D: 256 \times 128 \times 256 (16 \times 10 \times 16 cm)
 - Initial velocity conditions
 - 2D: 1024² (32 \times 32 cm) (in progress)
 - 3D: 256 \times 128 \times 256 (16 \times 10 \times 32 cm) (in progress)
- Simulation parameters chosen to match experimental conditions
 - ρ_1 = 0.9986 g/cm³, μ_1 = 0.009 g/(cm s)
 - ρ_2 = 0.9970 g/cm³, μ_2 = 0.011 g/(cm s)
 - $A = (\rho_1 \rho_2)/(\rho_1 + \rho_2) = 7.5 \times 10^{-4}$
 - Sc = \sqrt{D} = 7.0, $\nu = (\mu_1 + \mu_2)/(\rho_1 + \rho_2)$
 - $g_z = -981 \text{ cm/s}^2$





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PLIF density images from experiment are qualitatively similar to both 2D and 3D DNS densities



| PLIF Image [†] | 2D DNS | 3D DNS |
|-------------------------|-----------------------|-----------|
| <i>τ</i> = 1.15–1.30 | <i>τ</i> = 2.1 | au = 1.95 |

- PLIF images closely resemble DNS, but more resolution is required in 3D to fully resolve all scales present in experiment at later times
- Time lag between experiment and simulations with initial density conditions noted with respect to times of similar development

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3D DNS visualizations show initially 2D behavior with 3D structure emerging at later times



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The evolution of the mixing layer width from 2D and 3D DNS show differences, but need to evolve to later times

- Slopes of front growths are slightly steeper in 3D than in 2D
- Unclear whether 3D simulations exhibit a τ^2 growth at late times



The evolution of α_b and outer-scale Reynolds number from 2D and 3D simulations are quite different

- $\alpha_b = h_b/(A g t^2)$ from DNS bracket experimental values 0.06-0.07
 - For present initialization, $\alpha_{3D} > \alpha_{2D}$, which differs from trend observed using monotone integrated large-eddy simulation (MILES)
 - Both simulations seeded with long wavelength perturbations, yet exhibit different late-time α_b
- Outer-scale Reynolds number $Re_h = (h/v) dh/dt \approx 3,500$ at late times



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The time-evolution of turbulent molecular mixing parameters from 2D and 3D simulations are similar

- θ quantifies degree of molecular mixing ($\theta = 1$ and 0 represent <u>completely</u> <u>mixed fluids</u> and <u>completely segregated fluids</u>, respectively)
- B_0 and B_2 quantify ρ fluctuations for a miscible and immiscible mixture
- Values of θ qualitatively agree with experiment, but <u>lag in time</u>
- Inclusion of initial velocity data expected to reduce/eliminate time-lag and improve agreement with experimental data

Density spectra from 2D and 3D DNS show generally good agreement for intermediate *k* at later times

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These are the first DNS of the TAMU Water Channel experiment using experimentally-measured data

- Similarly-initialized 2D and 3D simulations exhibit different mixing layer growth rates and α_b
- Mixing parameters from both 2D and 3D simulations are in good agreement with measured values, but exhibit early-time lag
- Late-time density spectra are in generally good agreement with measured spectra, and also exhibit early-time lag
- Addition of initial velocity in DNS may be needed to match experimental data
- Simulations with *both* initial density and velocity conditions are in progress