



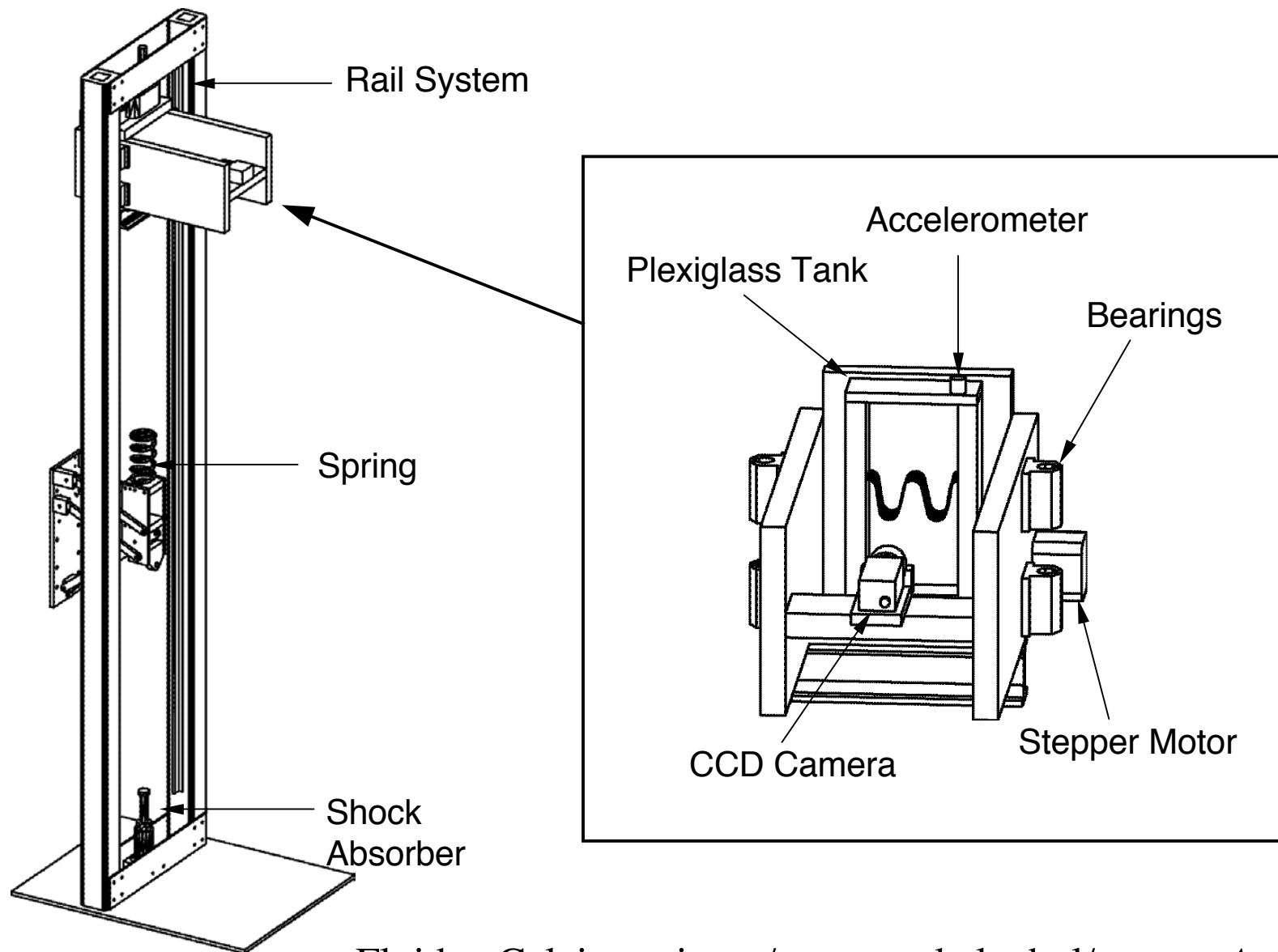
Experiments on the Single-Mode Three-Dimensional Rayleigh-Taylor and Richtmyer-Meshkov Instabilities

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Supported by LLNL and NASA

Richtmyer-Meshkov Apparatus

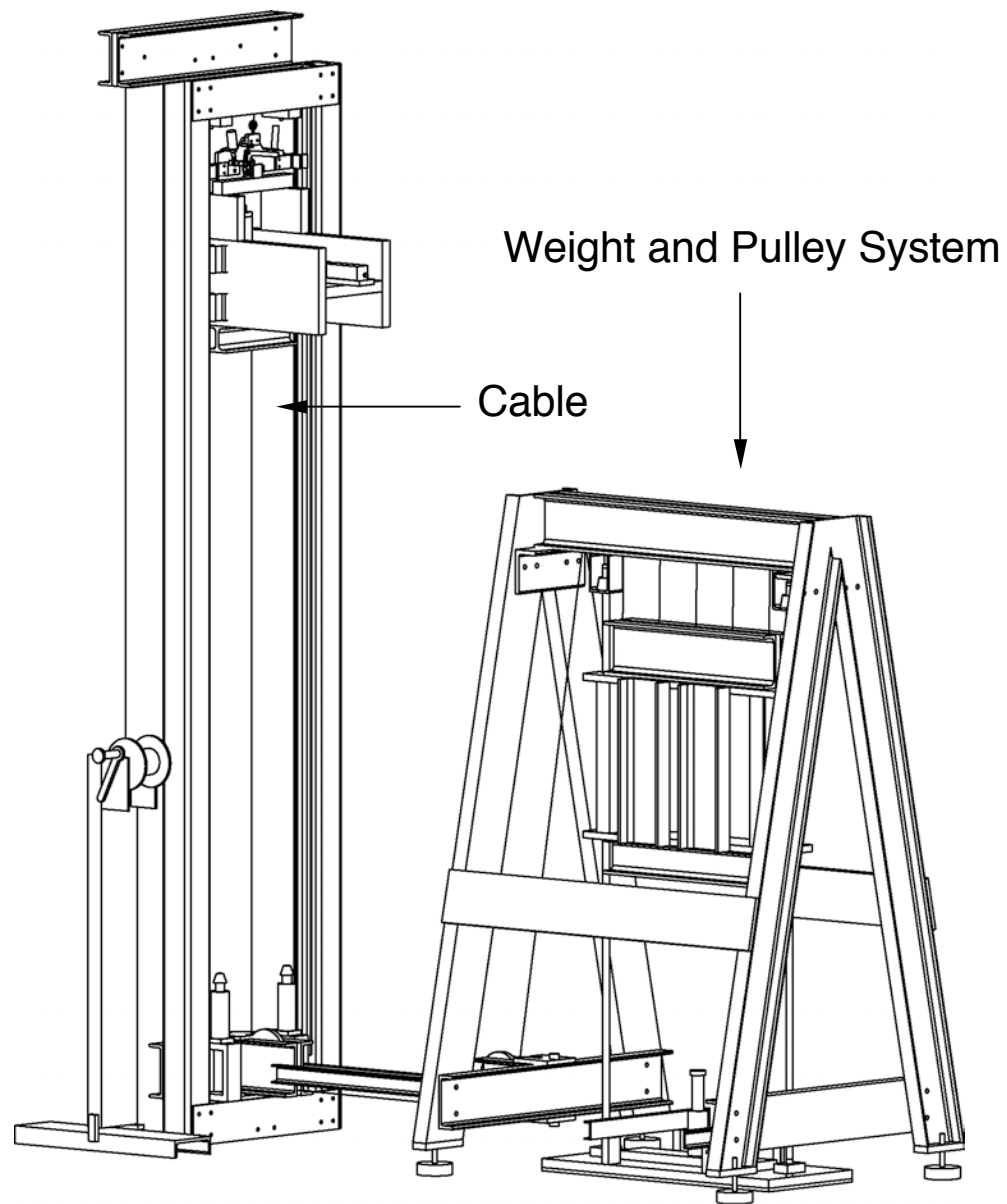


Fluids: Calcium nitrate/water and alcohol/water, $A \approx 0.15$

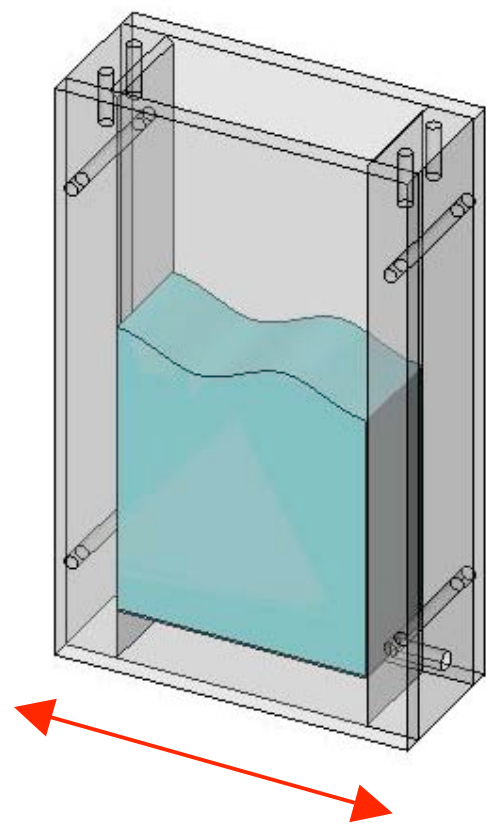
Drop Tower Operation



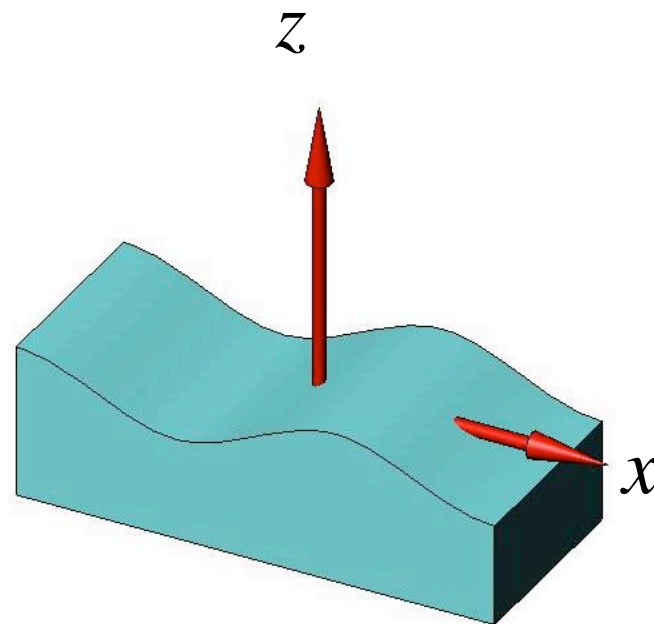
Rayleigh-Taylor Apparatus



2-D Perturbation

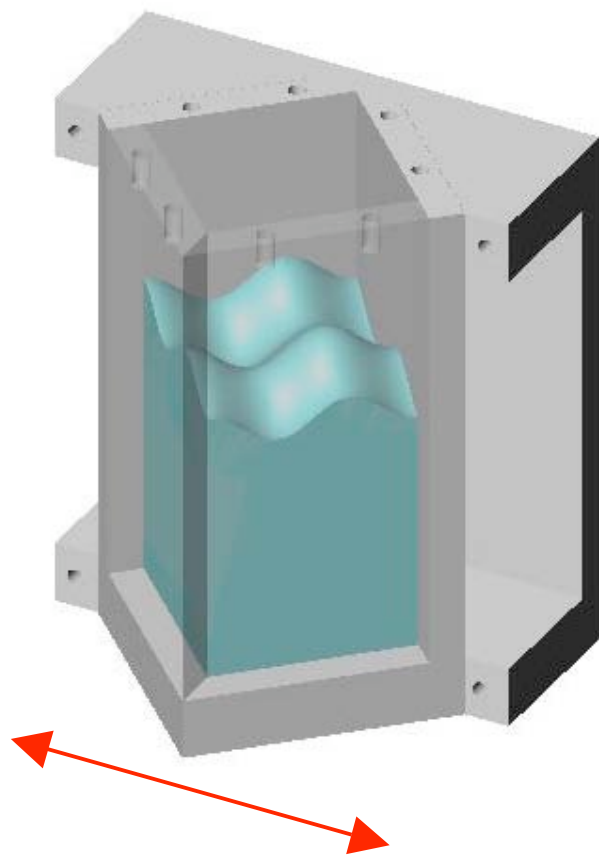


Direction of oscillation

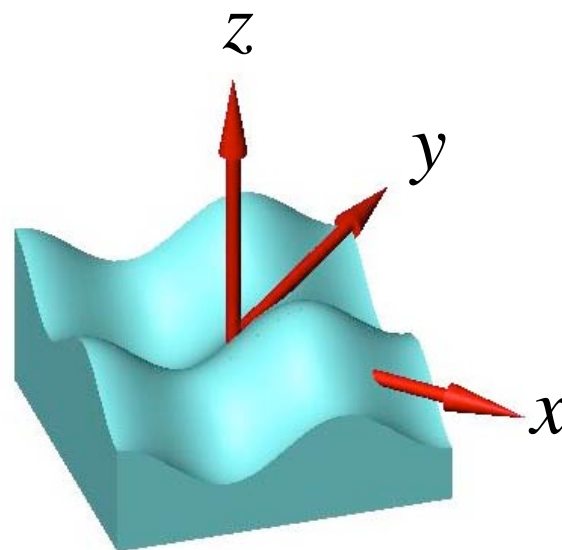


$$\eta = a_0 \sin(kx)$$

3-D Perturbation

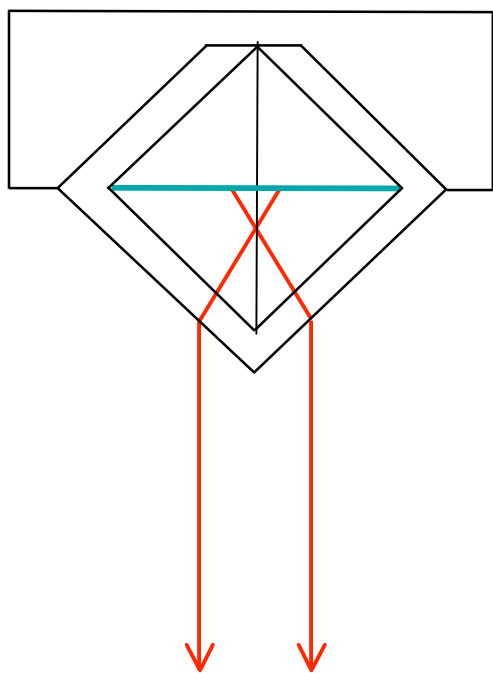


Direction of oscillation

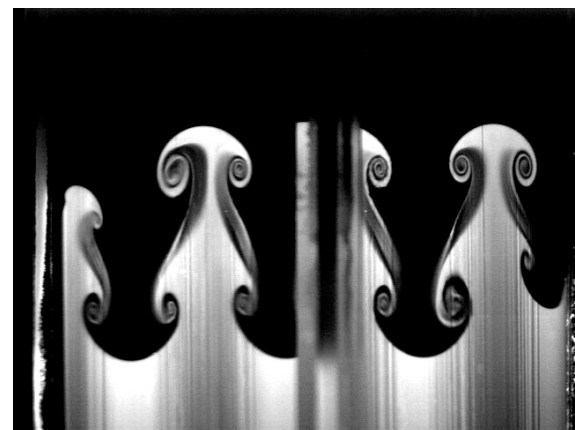


$$\eta = a_0 \sin(kx) \sin(ky)$$

Top view of tank



Refraction index differences causes light rays to be bent thus forming a double image.

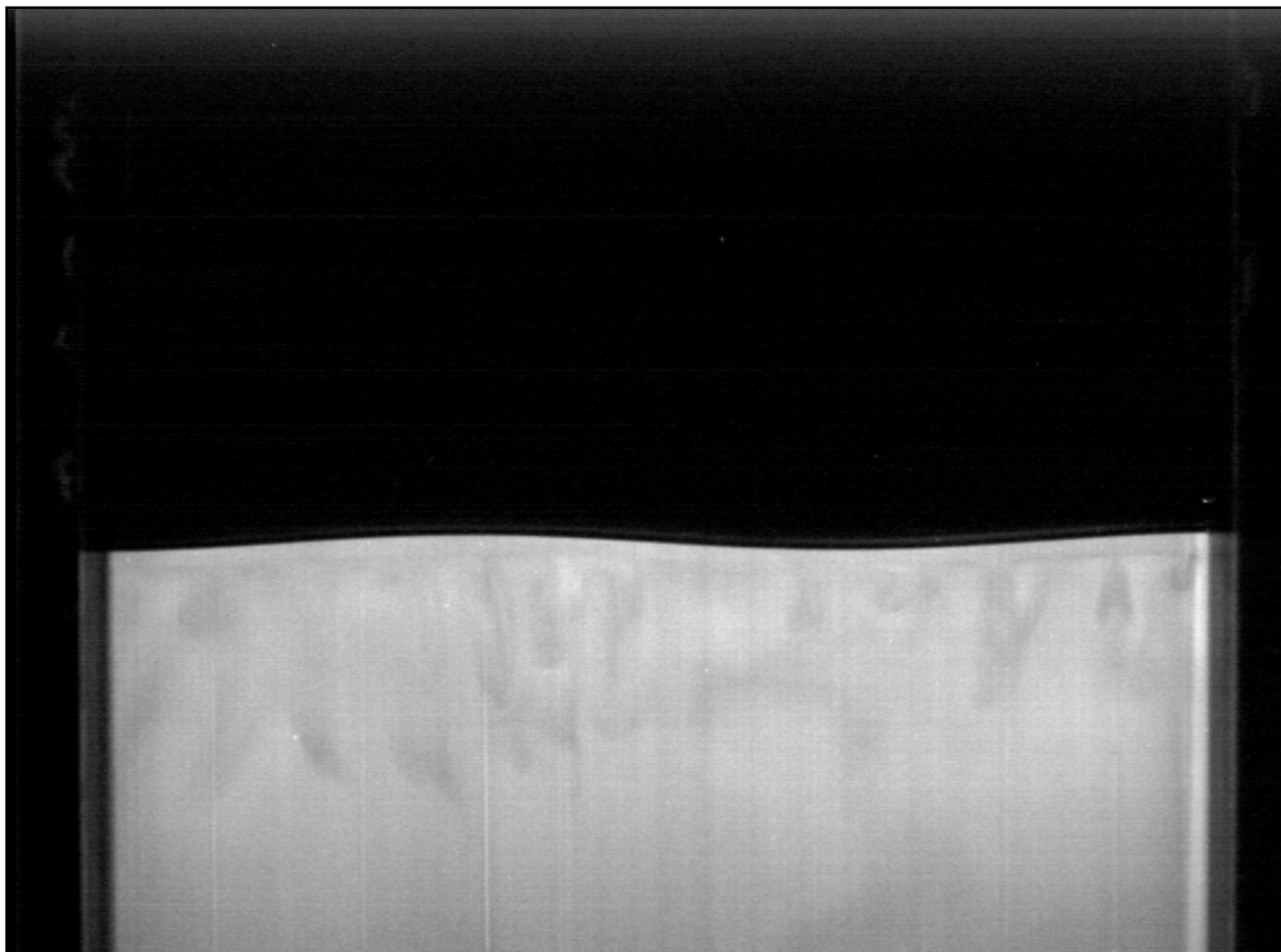


Original image

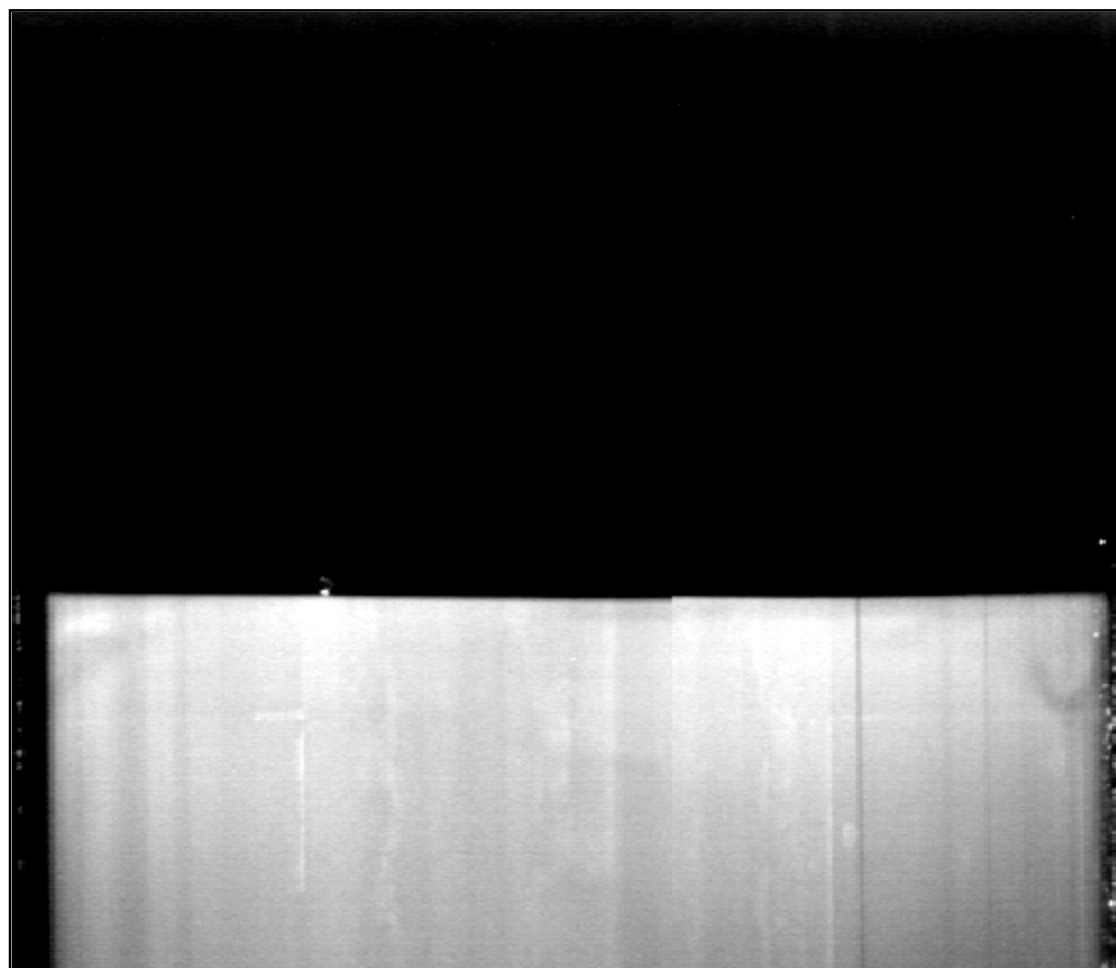


Corrected Image

2-D RM Instability



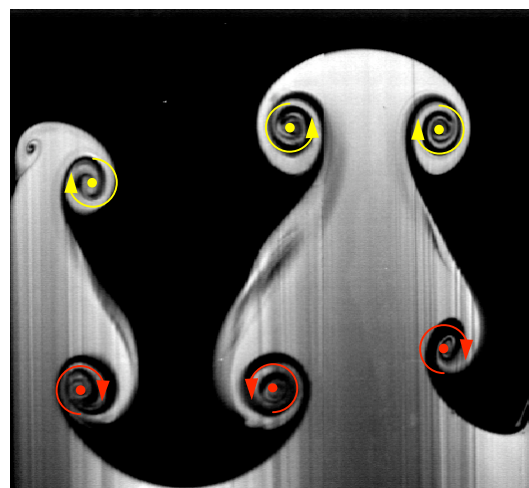
3-D RM Instability



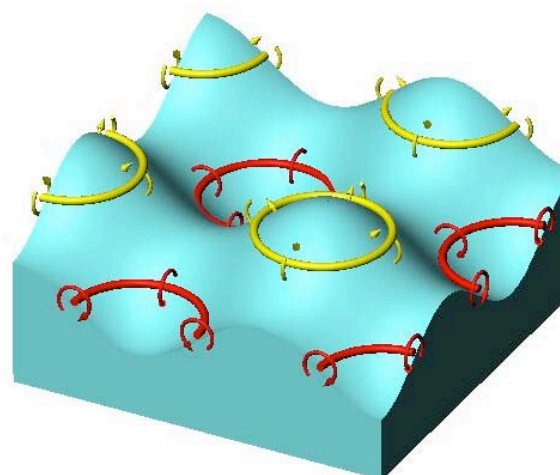
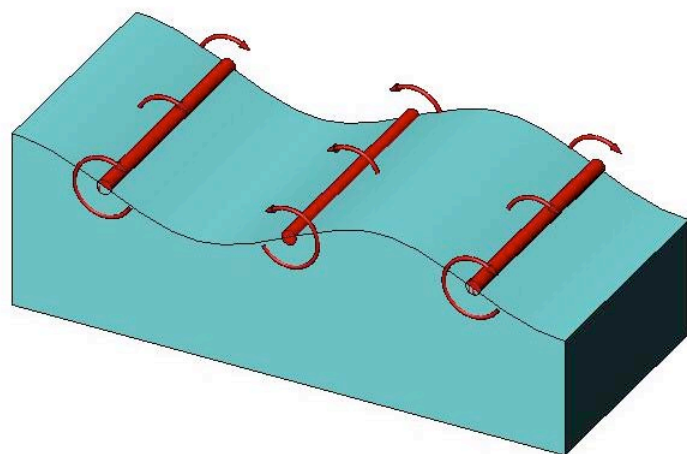
Late-Time Vorticity Distribution



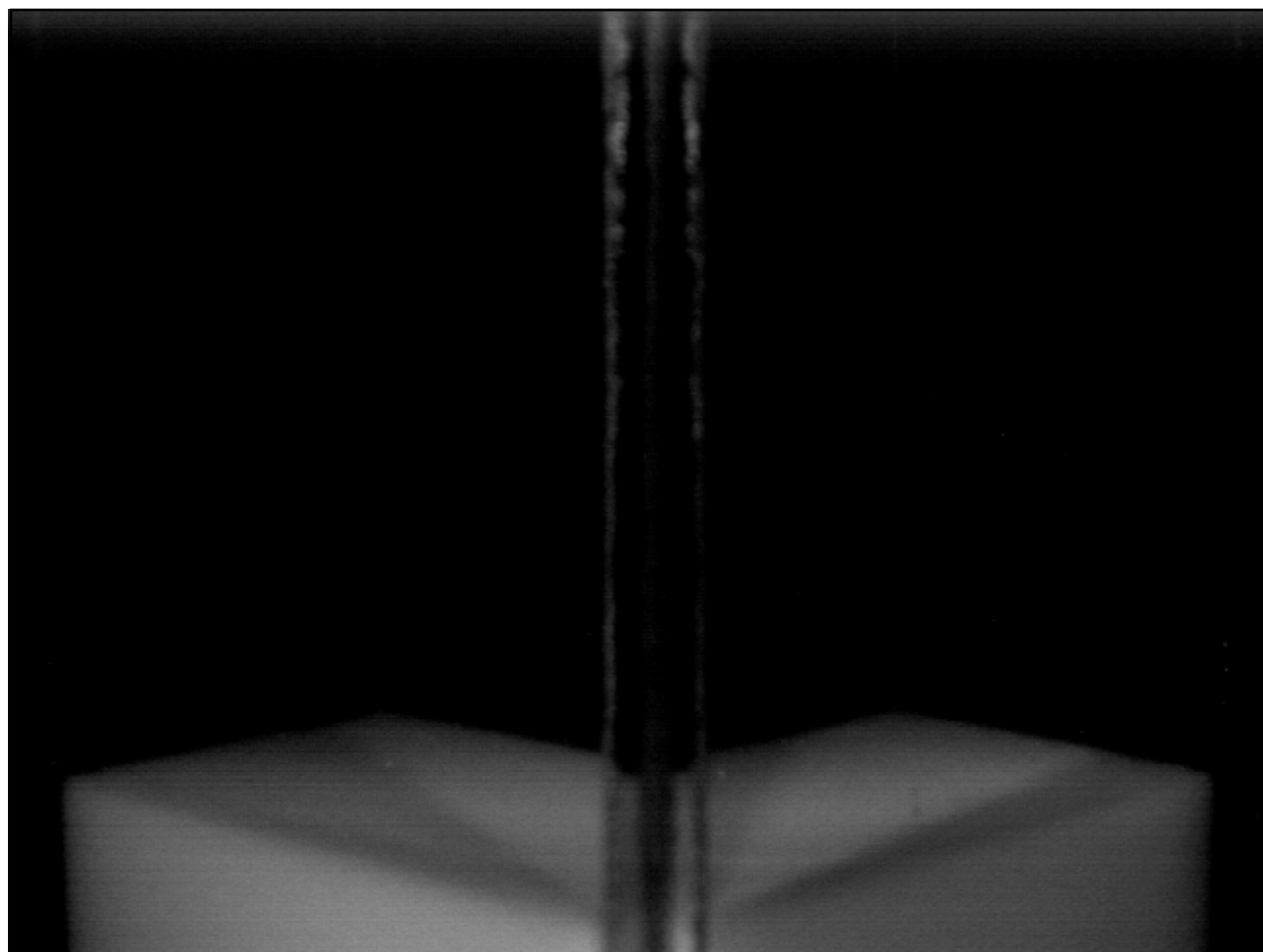
2-D



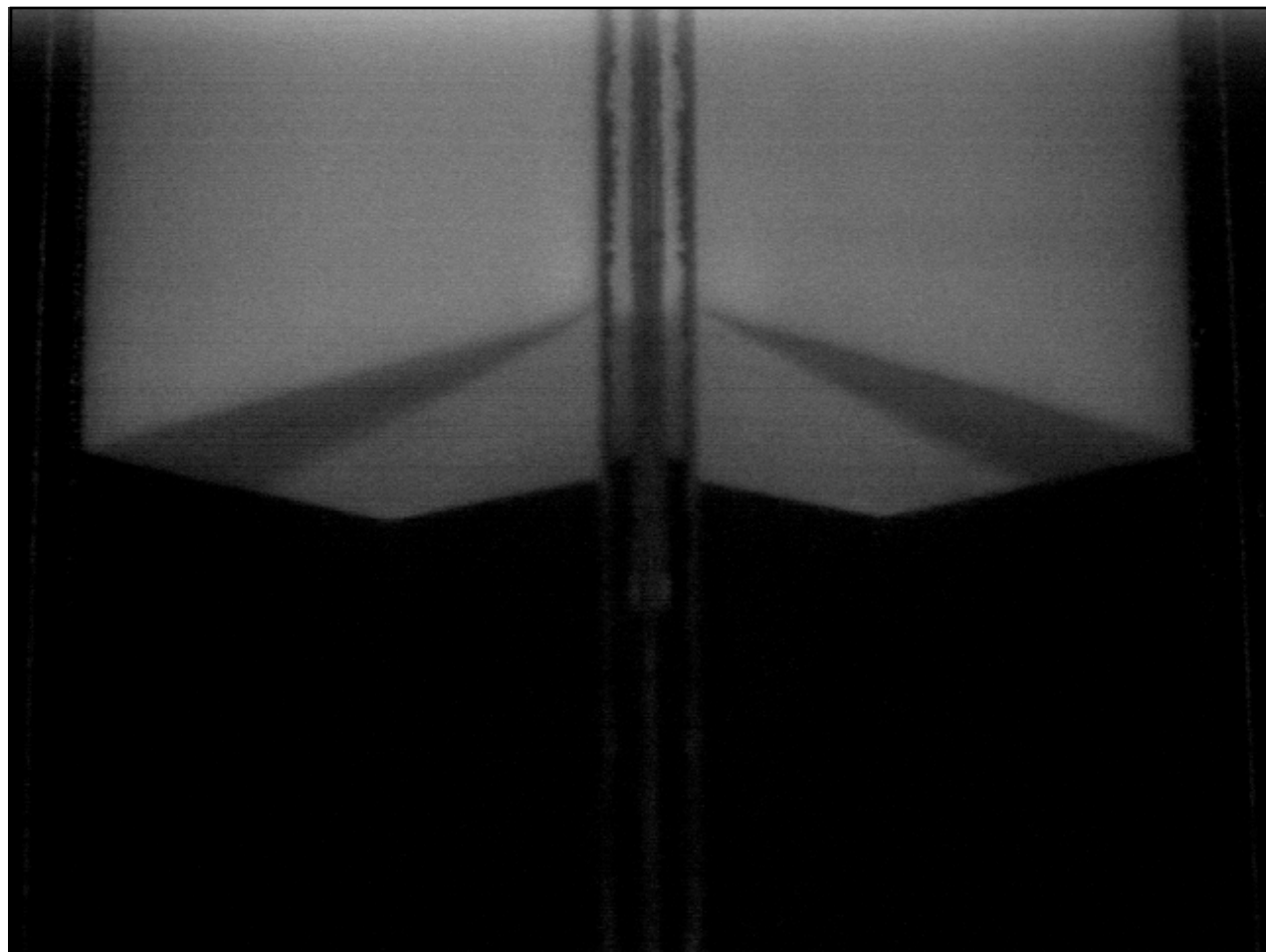
3-D



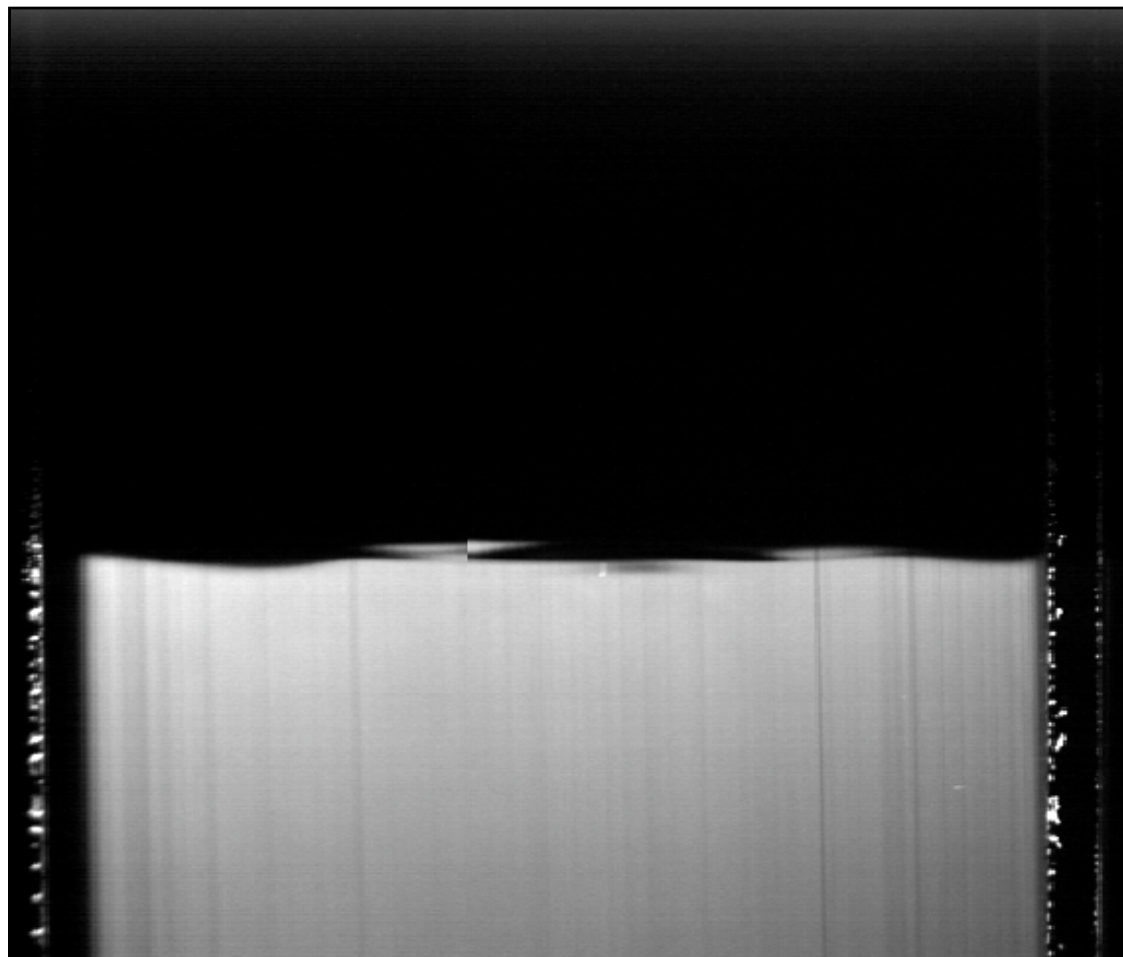
3-D RM Instability Opaque View



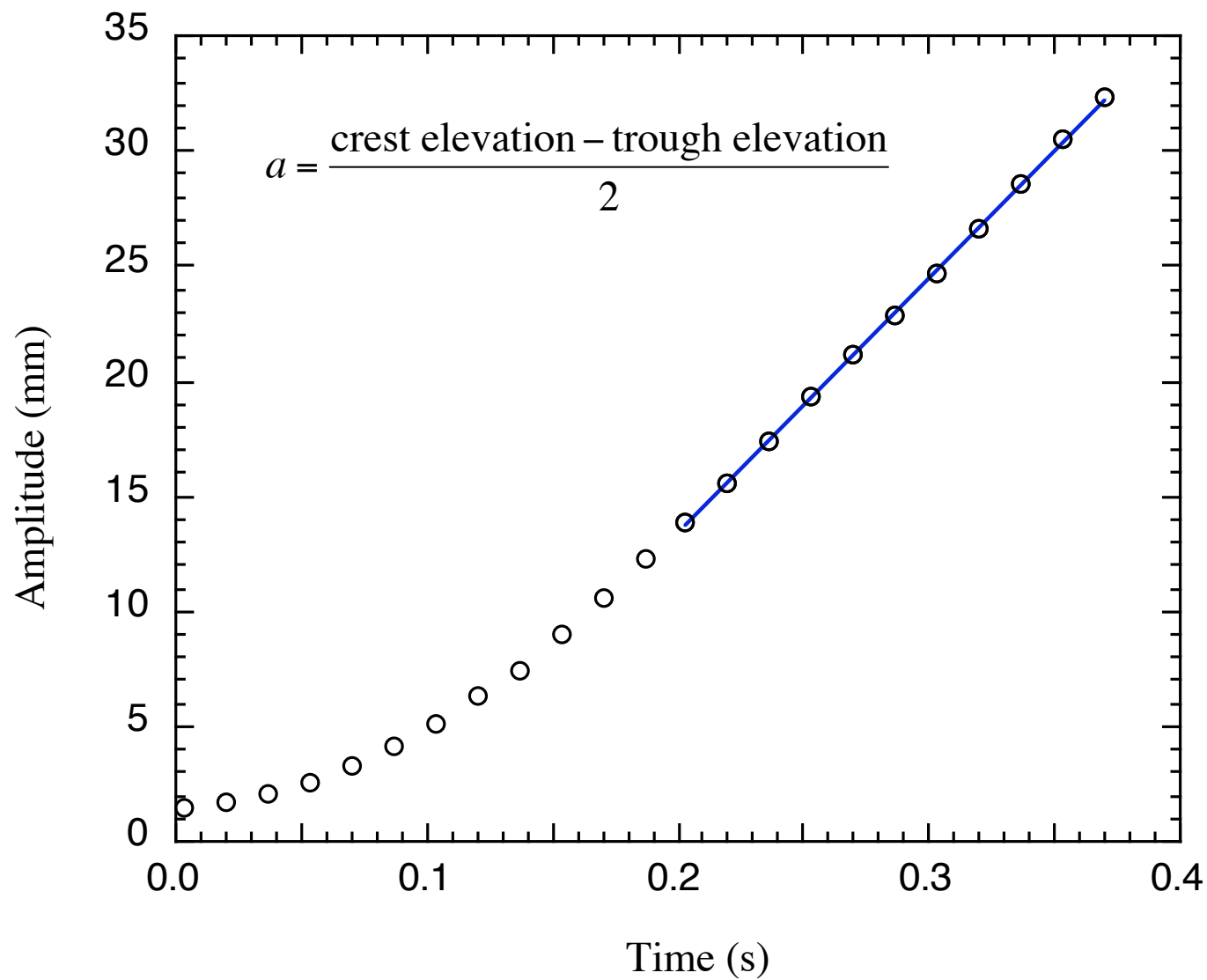
Opaque Bottom View



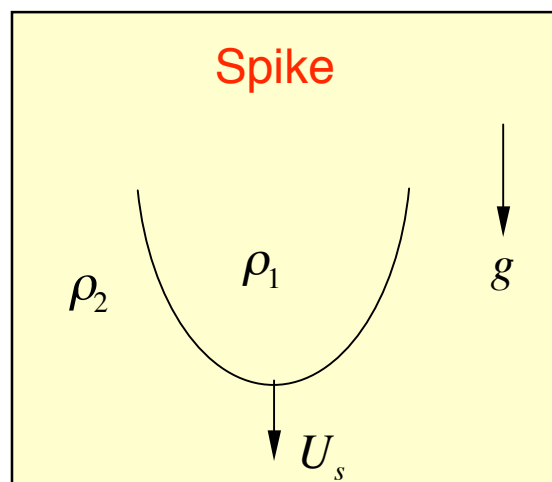
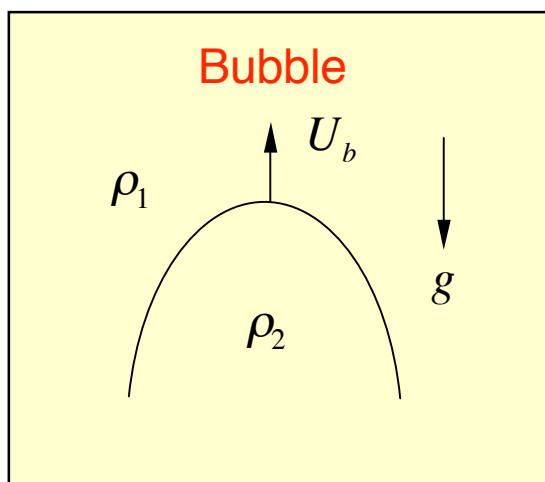
3-D RT Instability



RT Amplitude Measurements



Buoyancy/Drag Model



$$(\rho_2 + \kappa \rho_1)V\dot{U}_b = -C\rho_1 U_b^2 S + (\rho_1 - \rho_2)Vg$$

$$(\rho_1 + \kappa \rho_2)V\dot{U}_s = -C\rho_2 U_s^2 S + (\rho_1 - \rho_2)Vg$$

For RT instability

$$\dot{U}_{b/s} = 0$$

$$U_{b/s}^2 = \frac{Vk}{CS} \left(\frac{2A}{1 \pm A} \right) \frac{g}{k}$$

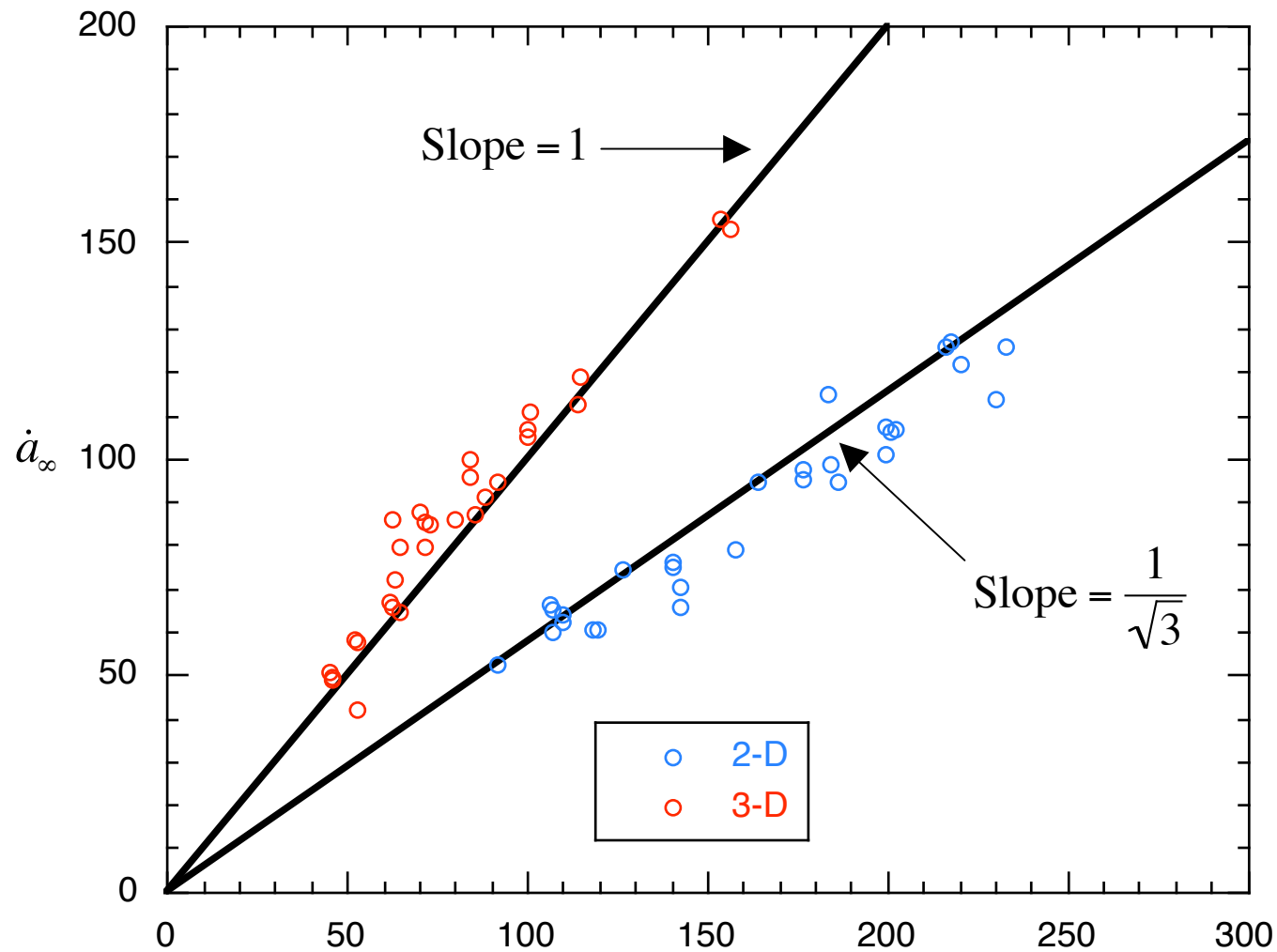
$$\frac{Vk}{CS} = \begin{cases} 1 & \text{for 2D} \\ 1/3 & \text{for 3D} \end{cases}$$

$$U_{b/s} = \begin{cases} \sqrt{\left(\frac{2A}{1 \pm A} \right) \frac{g}{3k}} & \text{for 2D} \\ \sqrt{\left(\frac{2A}{1 \pm A} \right) \frac{g}{k}} & \text{for 3D} \end{cases}$$

Obtained by Oron et al. (2001), Goncharov (2002)

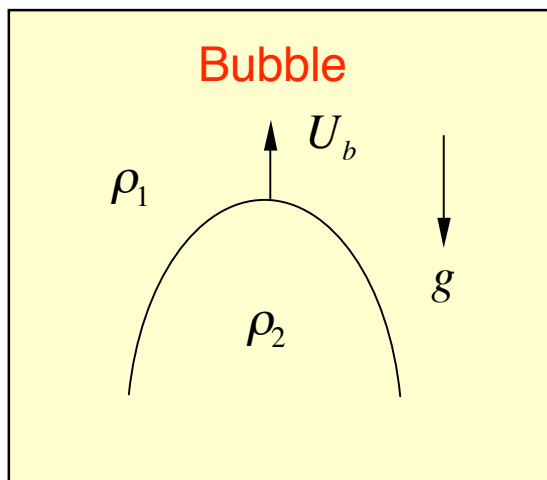
Also in agreement with Layzer (1955) when $A \rightarrow 1$

RT Late-Time Velocity



$$\frac{1}{2} \left[\sqrt{\left(\frac{2A}{1+A}\right) \frac{g}{k}} + \sqrt{\left(\frac{2A}{1-A}\right) \frac{g}{k}} \right]$$

Buoyancy/Drag Model



$$(\rho_2 + \kappa \rho_1)V\dot{U}_b = -C\rho_1 U_b^2 S + (\rho_1 - \rho_2)Vg$$

For RM instability $g = 0$

$$U_{b/s}^2 = -\frac{Vk}{CS} \left(\frac{1 \mp A}{1 \pm A} + \kappa \right) \frac{\dot{U}_{b/s}}{k}$$

As before:

$$\frac{Vk}{CS} = \begin{cases} 1/3 & \text{for 2D} \\ 1 & \text{for 3D} \end{cases}$$

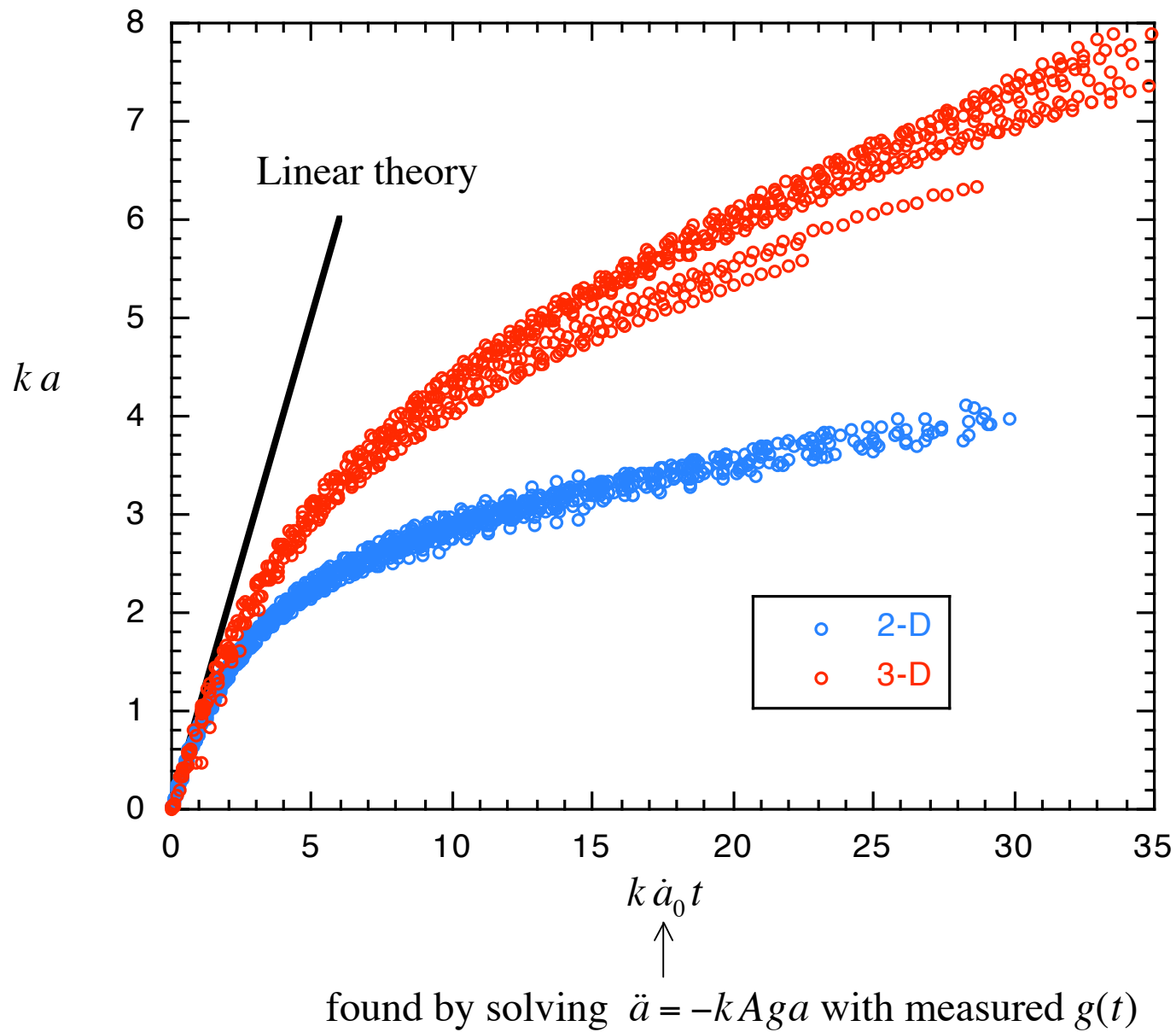
New parameter:

$$\kappa = \begin{cases} 1 & \text{for 2D} \\ 2 & \text{for 3D} \end{cases}$$

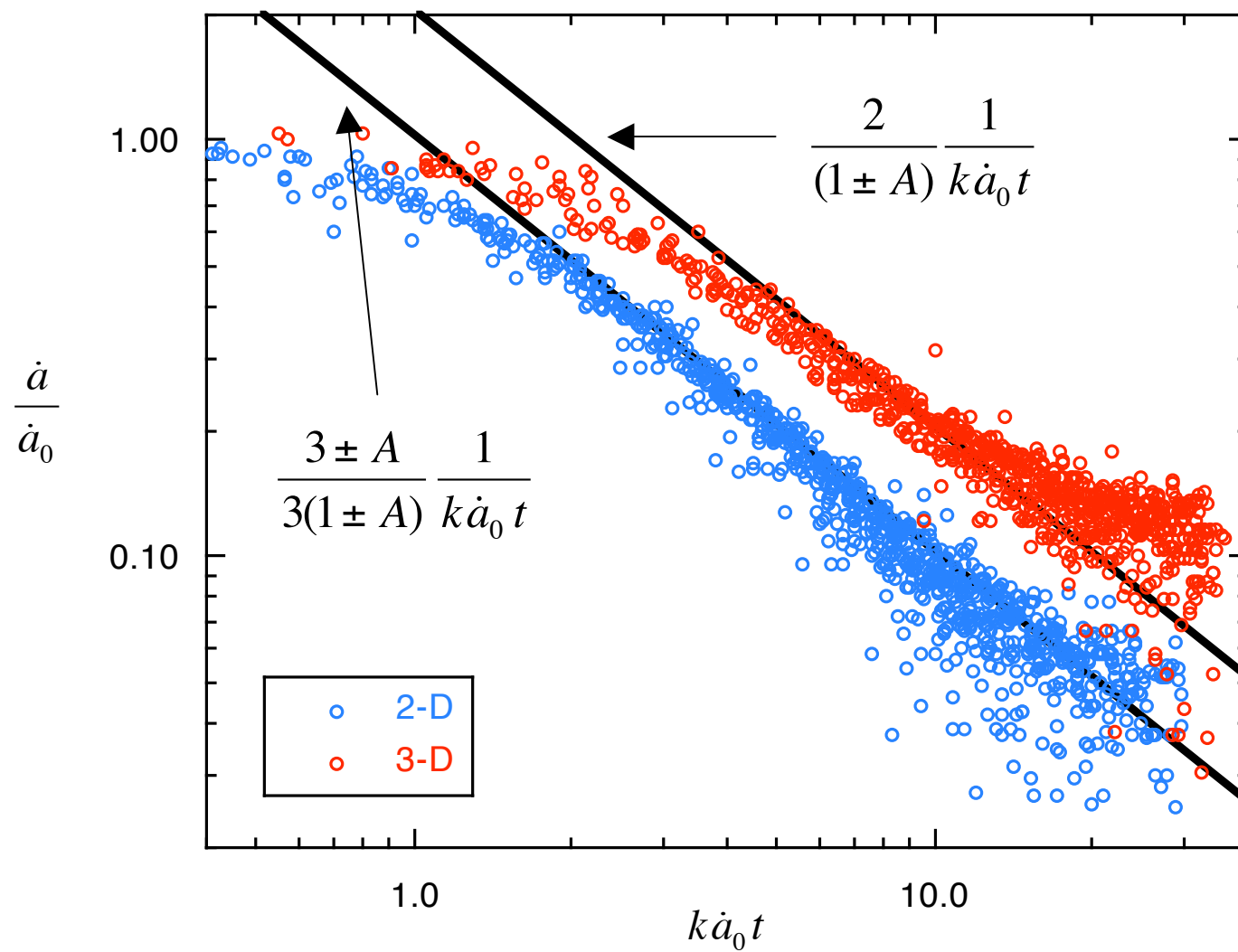
$$U_{b/s} = \begin{cases} \frac{3 \pm A}{3(1 \pm A)} \frac{1}{kt} & \text{for 2D} \\ \frac{2}{(1 \pm A)} \frac{1}{kt} & \text{for 3D} \end{cases}$$

In agreement with vortex model when $A \rightarrow 0$

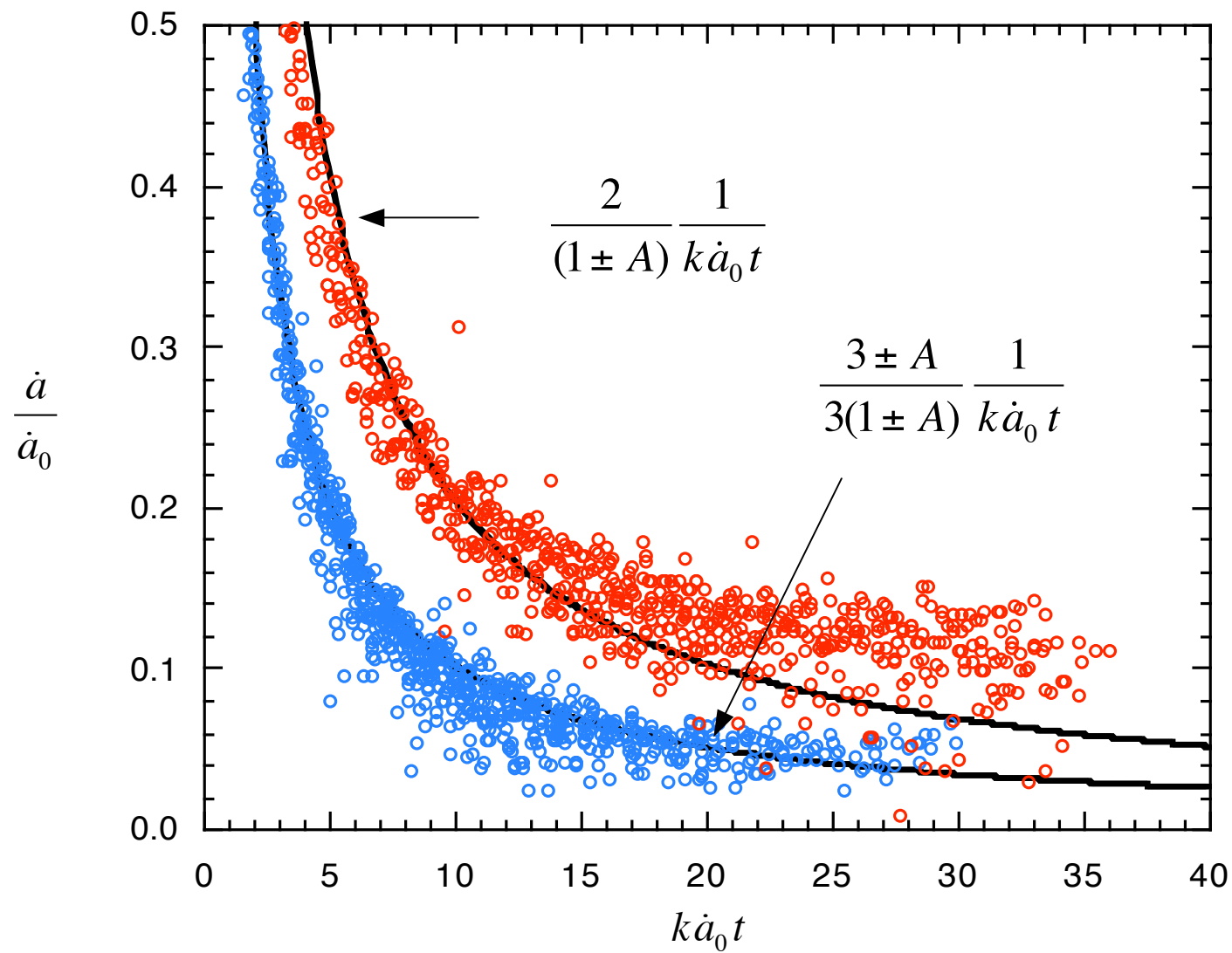
RM Amplitude Measurements



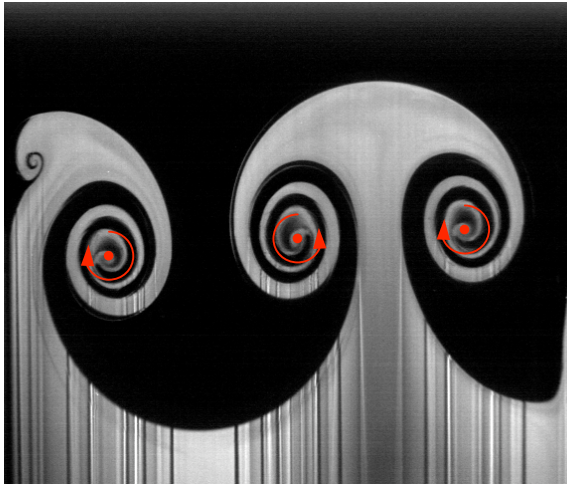
RM Late-Time Velocity



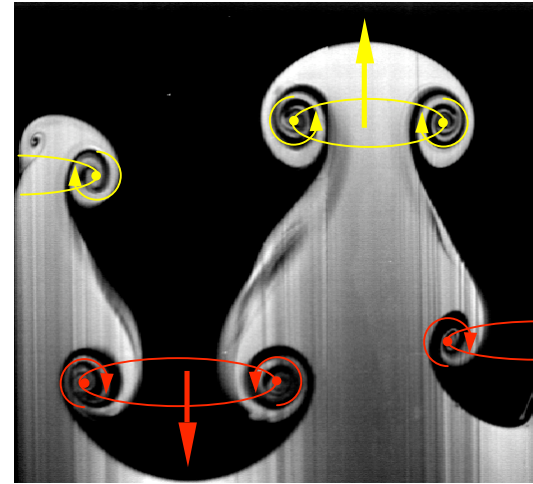
RM Late-Time Velocity



Vortex Dynamics



- The 2-D perturbation generates a nearly stationary vortex row.
- The result is bubble and spike velocities that decay as $1/kt$.



- The 3-D generates a square array of vortex rings that self-induce motion upward and downward.
- The result is a bubble and spike amplitude with constant velocity.

Low Atwood number RM instability is a vortex dominated.

Thus bubble models are not well suited for describing this flow.

Conclusions

- Low Atwood number three-dimensional single-mode RM and RT instabilities look fundamentally different than their two-dimensional counterparts.
- Late-time 3D RT amplitude measurements shows faster growth than 2D in agreement with the standard models.
- Late time 3D RM amplitude measurements also show significantly faster growth but appear to diverge from the predicted $1/kt$ dependence as $t \rightarrow \infty$.
- The disagreement with the models can be attributed to the fact that bubble models are ineffective in modeling the vortex dynamics of this flow.