

Combined Shear
and
Buoyancy Instabilities

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Spectral Turbulence Model for Variable Density Turbulence

- Steinkamp, Clark and Harlow, Int. J. of Multiphase Flow (1999)

- $$\hat{R}_{ij}(\vec{x}, \vec{k}, t) = \iiint_{-\infty \text{ to } \infty} R_{ij}(\vec{x}_1, \vec{x}_2, t) e^{-i\vec{k}\vec{r}} d\vec{r}$$

where $\vec{x} = \frac{\vec{x}_1 + \vec{x}_2}{2}$, $\vec{r} = \vec{x}_2 - \vec{x}_1$

- $$R_{ij}(\vec{x}, \vec{k}, t) = \int_{\Omega_k} \hat{R}_{ij}(\vec{x}, \vec{k}, t) \frac{k^2 d\Omega_k}{(2\pi)^3}$$

- $$R_{ij}(\vec{x}, t) = \int_0^{\infty} R_{ij}(\vec{x}, k, t) dk$$

Advantages of Spectral Transport Models Over Single-Point Models (i.e., k - ϵ):

- More generality, such as in the case of non-equilibrium transients
- Does not require a length scale or dissipation equation
- Greater flexibility with modeling, such as with non-local interactions in both physical and wavenumber space

Disadvantages:

- Greater Complexity
- More computationally expensive
- More Flexibility!

Governing Equations

$$\frac{D\tilde{u}_i}{Dt} = -\frac{1}{\bar{\rho}} \left[\frac{\partial R_{in}}{\partial x_n} + \frac{\partial \bar{p}}{\partial x_i} \right] + v_m \frac{\partial^2 \tilde{u}_i}{\partial x_n^2} + g_i$$

$$\tilde{u}_i = \frac{\overline{\rho u_i}}{\bar{\rho}}, \quad R_{in} = \overline{\rho u_i'' u_n''} \quad \text{where } u_i = \tilde{u}_i + u_i''$$

$$\frac{\partial \bar{\rho}}{\partial t} + \frac{\partial \bar{\rho} \tilde{u}_n}{\partial x_n} = 0$$

$$\begin{aligned} \frac{DR_{ij}(k)}{Dt} = & \iiint_{-\infty t_0 \infty} \left[a_i(k) \frac{\partial \bar{p}}{\partial x_j} + a_j(k) \frac{\partial \bar{p}}{\partial x_i} \right]_{\bar{x}'} \hat{f}(\bar{x}, \bar{x}') d\bar{x}' - R_{in} \frac{\partial \tilde{u}_j}{\partial x_n} - R_{jn} \frac{\partial \tilde{u}_i}{\partial x_n} \\ & + c_d \frac{\partial}{\partial x_n} \left[v_t \frac{\partial R_{ij}}{\partial x_n} \right] + c_m \left(\frac{1}{3} \delta_{ij} R_{mn} - R_{ij} \right) \int_0^\infty \sqrt{\frac{k R_{mn}}{\bar{\rho}}} dk \\ & + \frac{\partial}{\partial k} \left\{ k^2 \sqrt{\frac{k R_{mn}}{\bar{\rho}}} \left[-c_1 R_{ij} + c_2 k \frac{\partial R_{ij}}{\partial k} \right] \right\} \\ & + \text{Non-Local}(T(k)) \end{aligned}$$

$$\text{where } a_i(\bar{x}_1, \bar{x}_2, t) = -\overline{u_i''(\bar{x}_1) \rho(\bar{x}_1) v(\bar{x}_2)} \quad \text{and } v_t = \int_0^\infty \sqrt{\frac{k R_{mn}}{\bar{\rho}}} \frac{dk}{k^2}$$

Governing Equations (cont.)

$$\begin{aligned} \frac{Da_i(k)}{Dt} = & \frac{b(k)}{\bar{\rho}} \frac{\partial \bar{\rho}}{\partial x_i} - \frac{R_{in}}{\bar{\rho}^2} \frac{\partial \bar{\rho}}{\partial x_n} + c_d \frac{\partial}{\partial x_n} \left[v_t \frac{\partial a_i}{\partial x_n} \right] \\ & - \left[c_{r1} k^2 \sqrt{a_n a_n} + c_{r2} k \sqrt{\frac{k R_{mn}}{\bar{\rho}}} \right] a_i \\ & + \frac{\partial}{\partial k} \left\{ k^2 \sqrt{\frac{k R_{mn}}{\bar{\rho}}} \left[-c_1 a_i + c_2 k \frac{\partial a_i}{\partial k} \right] \right\} \end{aligned}$$

where $b(\bar{x}_1, \bar{x}_2, t) = \overline{-\rho'(\bar{x}_1) v'(\bar{x}_2)}$

$$\begin{aligned} \frac{Db(k)}{Dt} = & u_n \frac{\partial b}{\partial x_n} + \frac{2\bar{\rho} - \rho_1 - \rho_2}{\rho_1 \rho_2} \frac{\partial \bar{\rho} a_n}{\partial x_n} + c_d \frac{\partial}{\partial x_n} \left[v_t \frac{\partial b}{\partial x_n} \right] \\ & - c_{fb} \left[\bar{v}^2 \frac{\partial \bar{\rho}/\bar{v}}{\partial x_n} \right] \frac{\partial k a_n}{\partial k} + \frac{\partial}{\partial k} \left\{ k^2 \sqrt{\frac{k R_{mn}}{\bar{\rho}}} \left[-c_1 b + c_2 k \frac{\partial b}{\partial k} \right] \right\} \\ & - 2v_t k^2 b \end{aligned}$$

Non-Local Behavior in Physical and Wavenumber Space

Physical Space:

$$\frac{DR_{mn}}{Dt} = 2 \int_{-\infty}^{\infty} a_2(y') \frac{\partial \bar{p}}{\partial y'} \hat{f}(y, y') dy' + \dots$$

$$\int_{-\infty}^{\infty} \hat{f}(y, y') dy = 1$$

$$\hat{f}(y, y') = \frac{\exp(-2k|y - y'|)}{\int_{-\infty}^{\infty} \exp(-2k|y - y''|) dy''}$$

Wavenumber Space: Kraichnan and Spiegel Model for Energy Transfer (1962)

$$T(k) = \int [S_a(k/p) - S_e(k/p)] dp$$

$$S_e(p/k) = \eta \sqrt{pE(p)} E(p) \left(\frac{p}{k}\right)^m g\left(\frac{p}{k}\right)$$

$$S_a(k/p) = S_e(p/k)$$

$$T(k) = -\Sigma(k) + \int \Sigma(p) \check{f}(k, p) dp \quad \text{where } \Sigma(k) = \eta [kE(k)]^{\frac{3}{2}}$$

$$\text{and } \check{f}(k, p) = p^{-1} \left(\frac{k}{p}\right)^m g\left(\frac{k}{p}\right)$$

Modeling Constants

$c_1, c_2, c_m, c_{r1}, c_{r2}, c_d$ also η and $g(x)$

$c_1/c_2 = 2$ for equipartition of energy ($c_1 = .5$)

$c_{r1} \approx c_{r2} = 5.0$ to match buoyancy experiments

$c_m = 1.0$ original SCH model constant

Given that $\int \tilde{f}(k, p) dk = 1 \Rightarrow \int x^m g(x) dx = 1$ where $x = k/p$

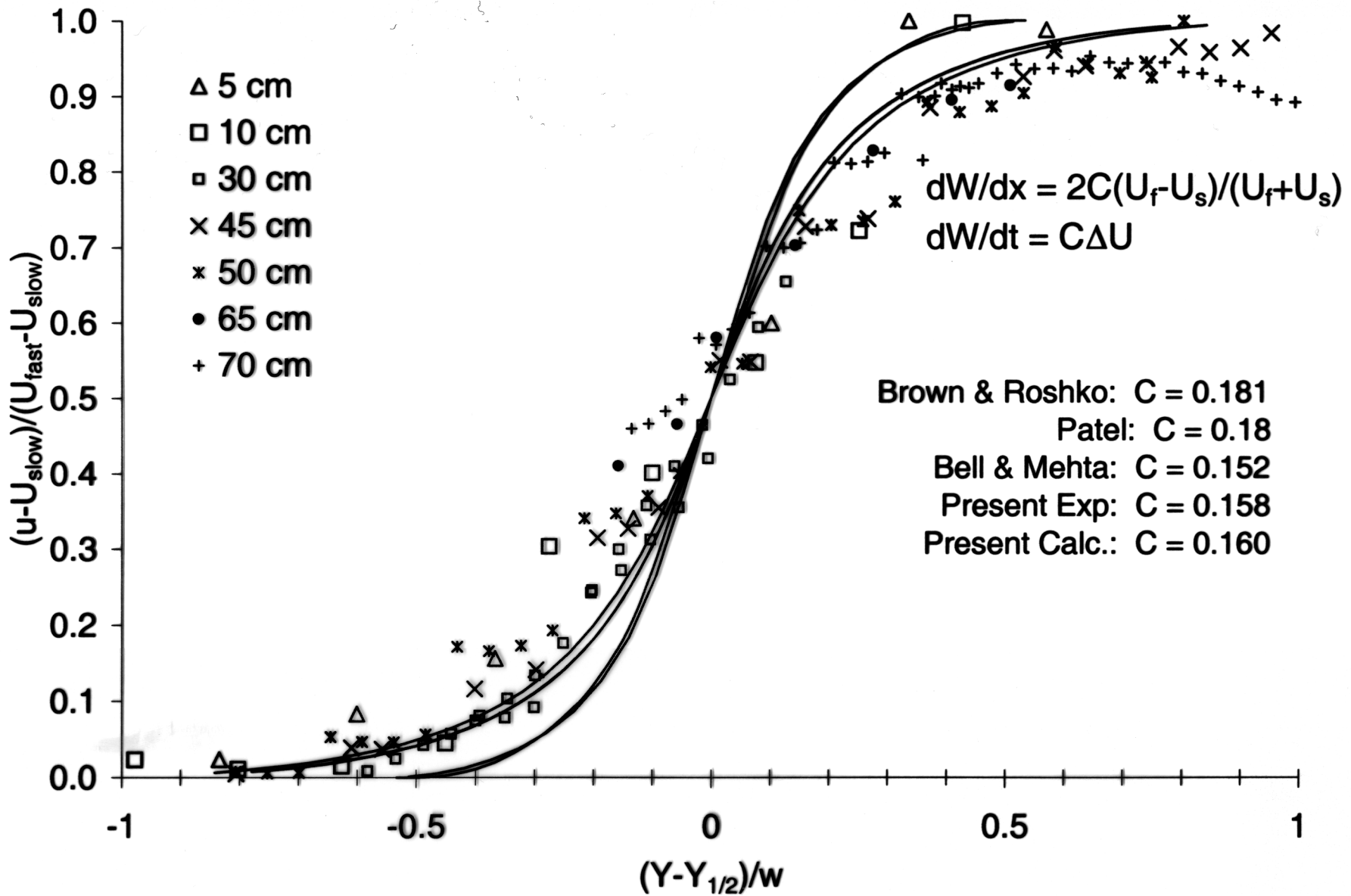
Let $g(x) = \exp[-\alpha(x + x^{-1})]/N$

and for $2k$ coupling ($m = -2, \alpha = 2/3$)

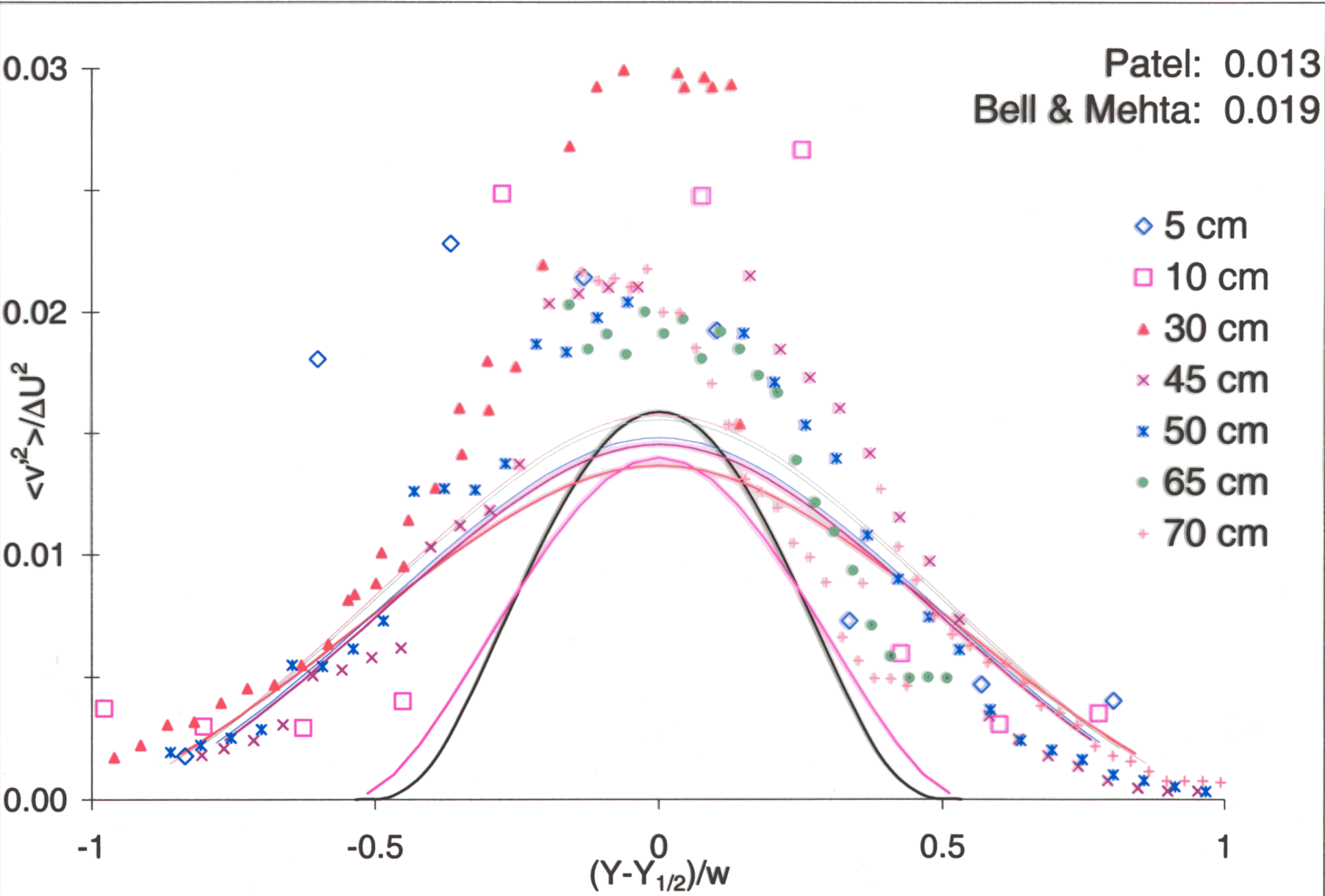
Lastly,

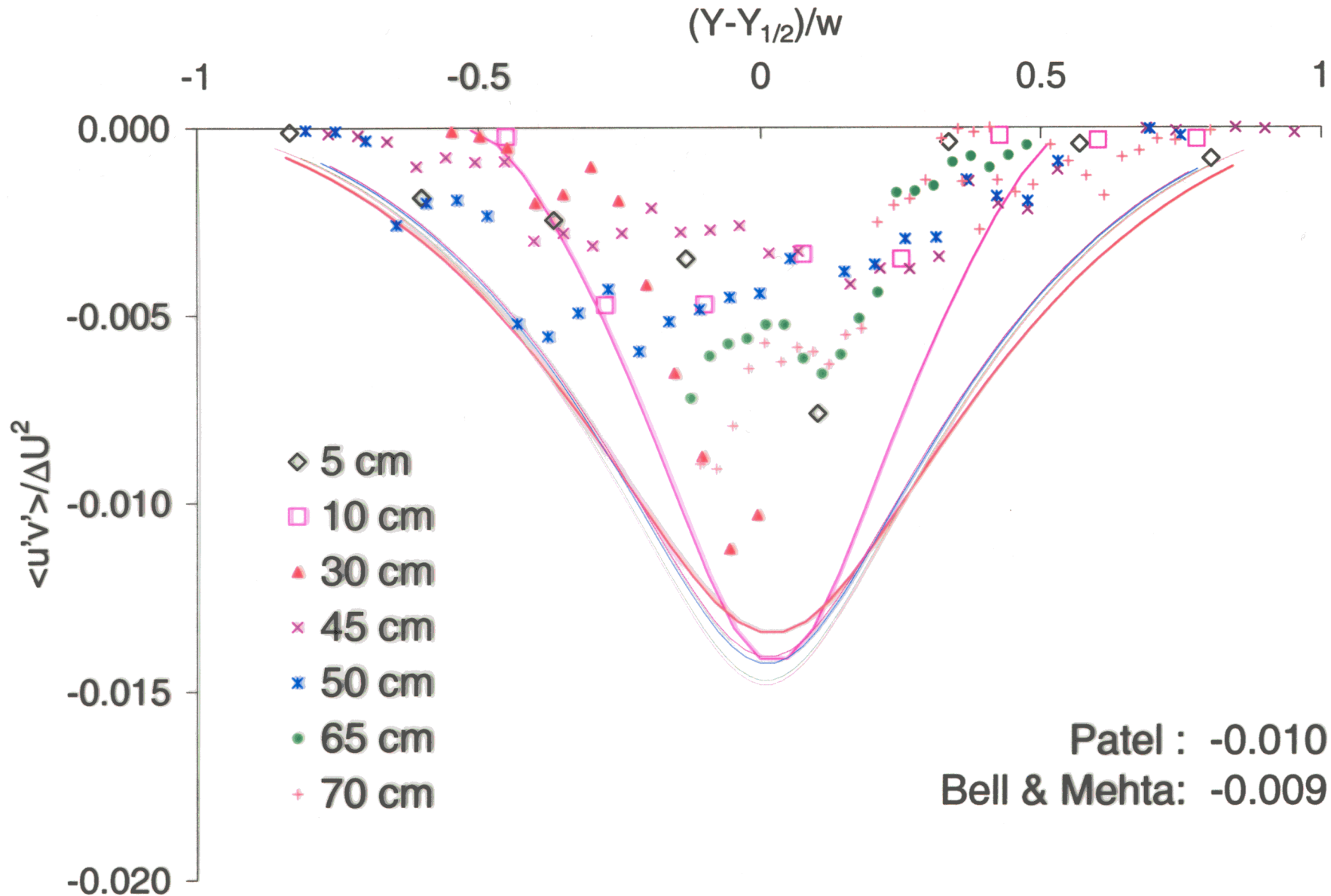
$$\varepsilon = \int_{k=0}^{\infty} \int S(q/p) dp dq$$

$$\eta = C_k^{-3/2} \left[\int_1^{\infty} \ln(x) (x^m - x^{-m-2}) g(x) dx \right]^{-1} \Rightarrow \eta = 1.36$$

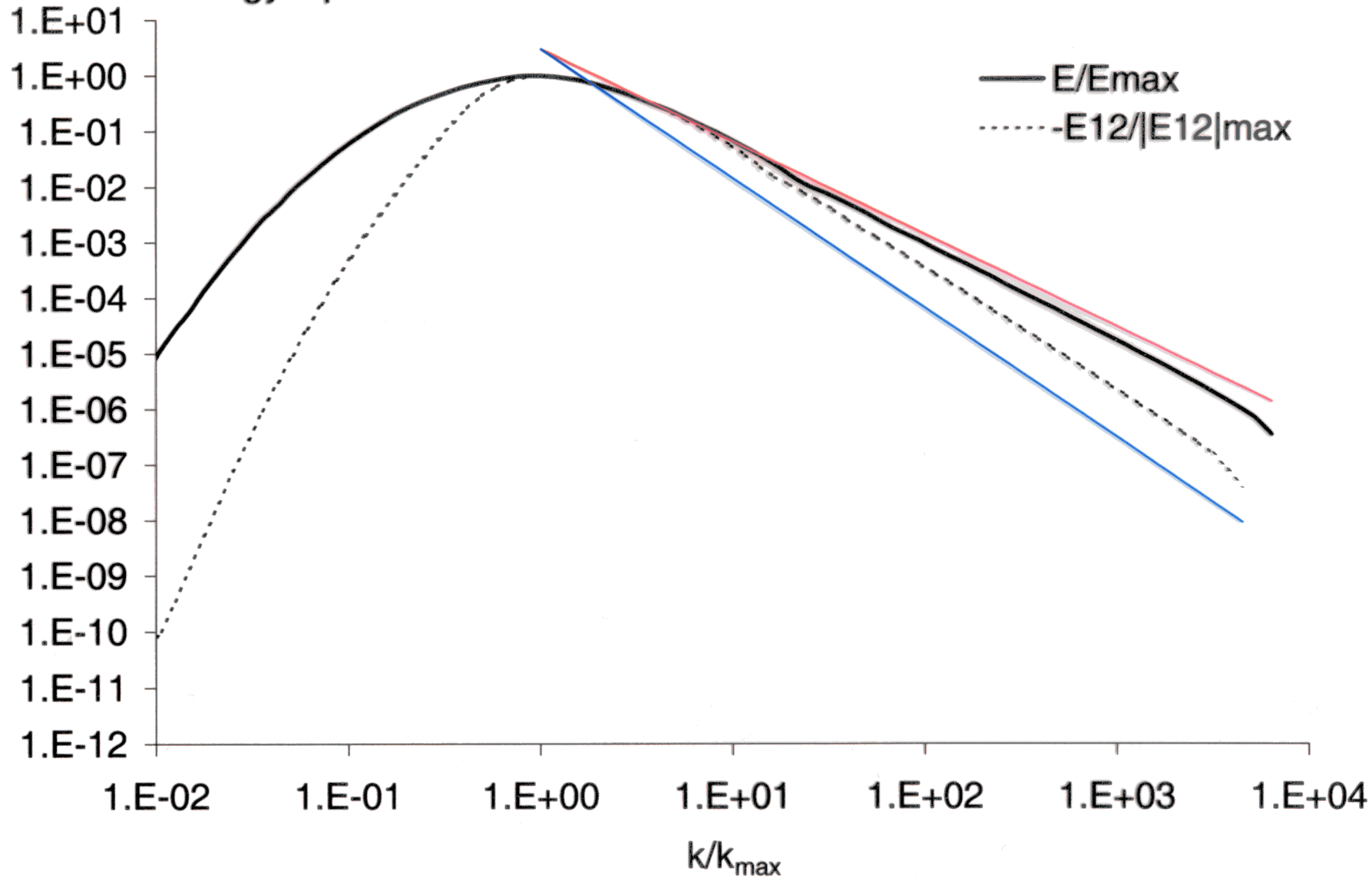


Patel: 0.013
Bell & Mehta: 0.019

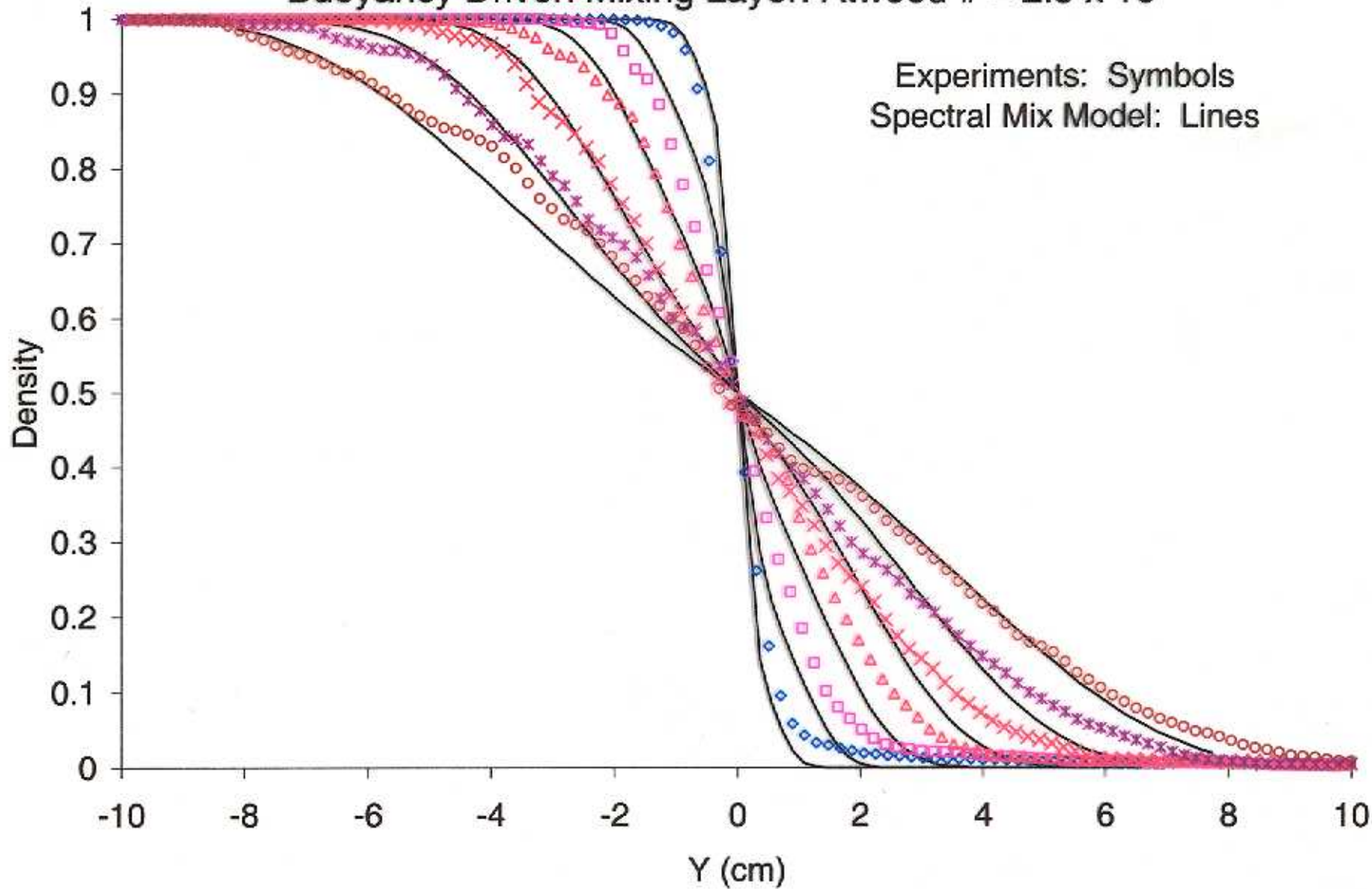




Shear Energy Spectra



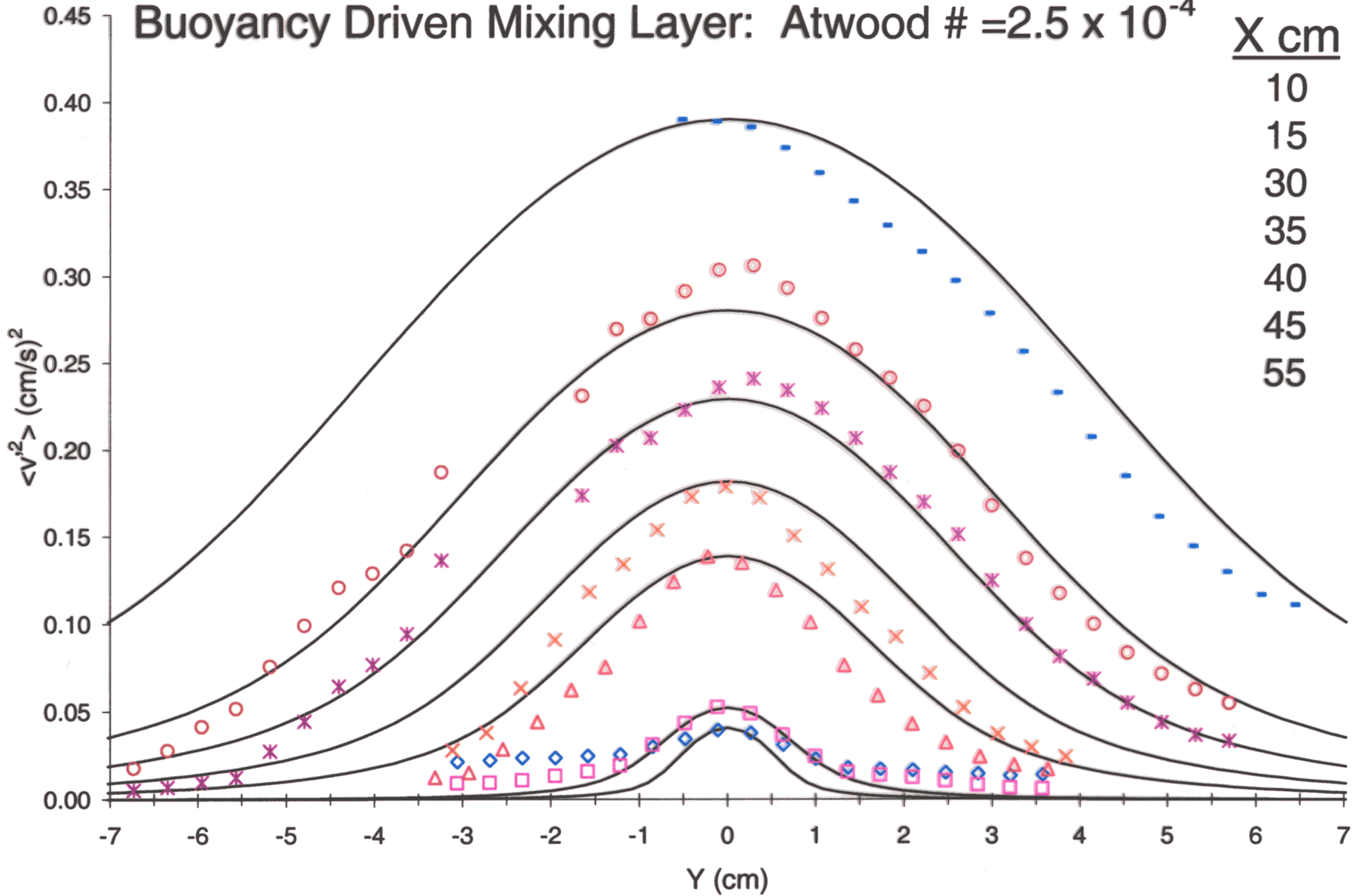
Buoyancy Driven Mixing Layer: Atwood # = 2.5×10^4



Buoyancy Driven Mixing Layer: Atwood # = 2.5×10^{-4}

X cm

- 10
- 15
- 30
- 35
- 40
- 45
- 55



b profiles in a buoyancy-driven mixing layer

