

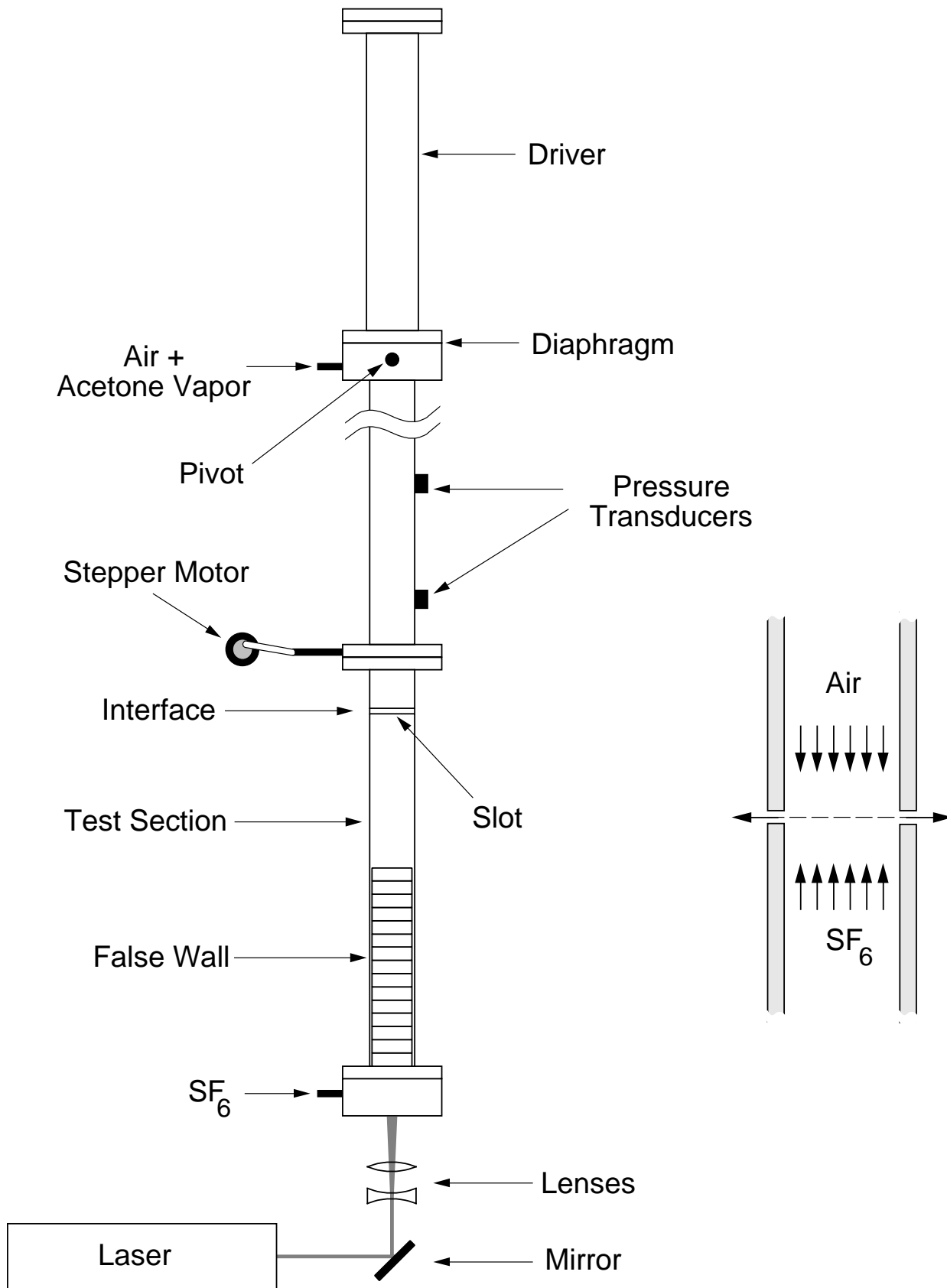
Experimental Investigation of Richtmyer-  
Meshkov Instability after a Second Interaction  
with a Reflected Shock Wave

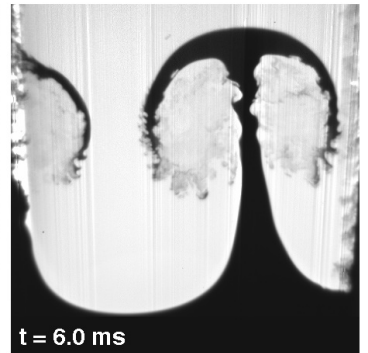
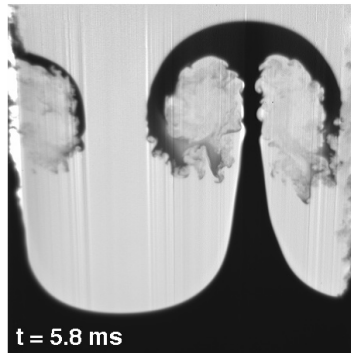
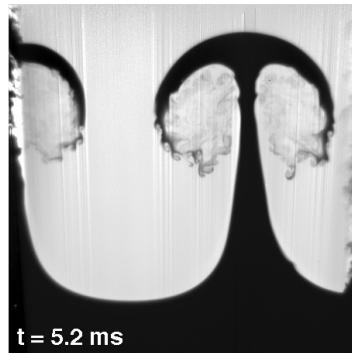
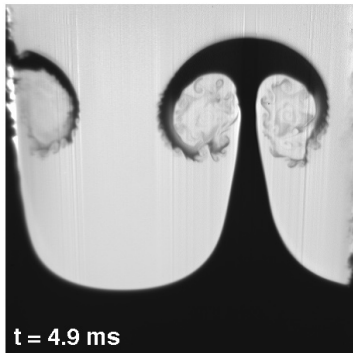
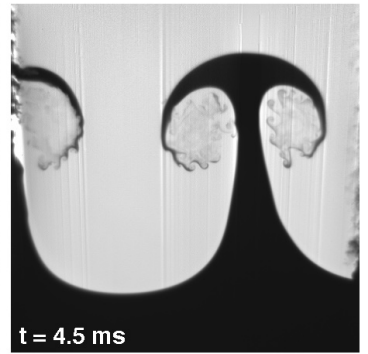
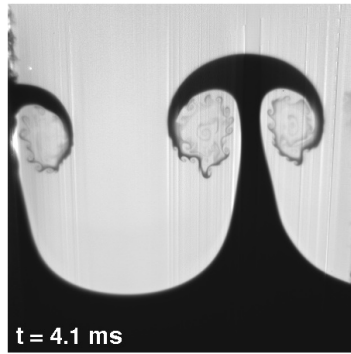
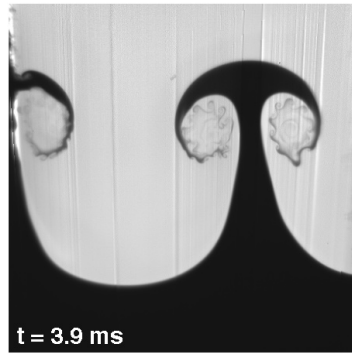
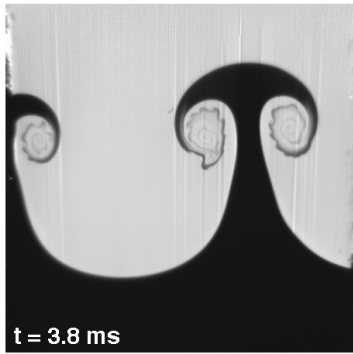
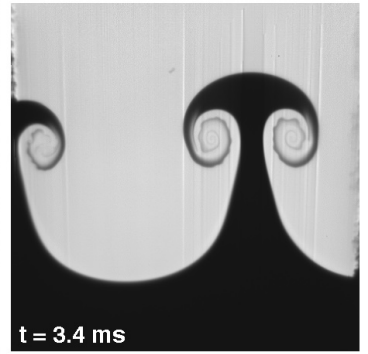
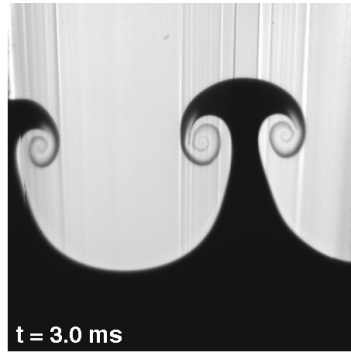
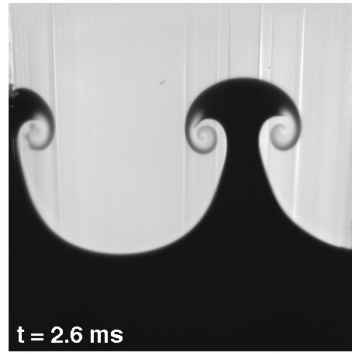
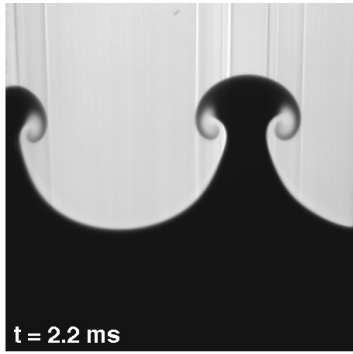
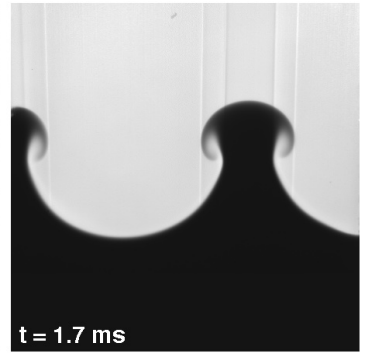
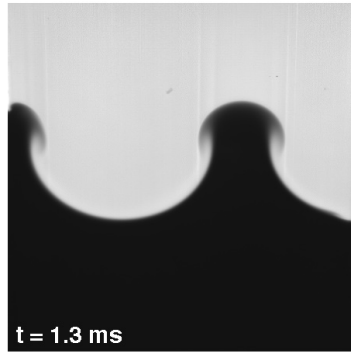
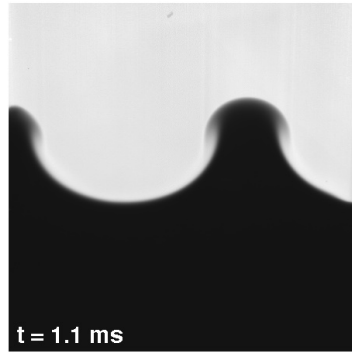
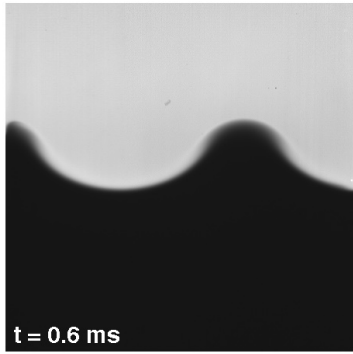
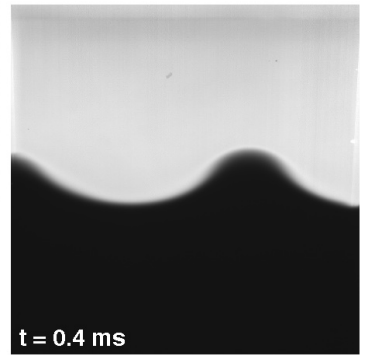
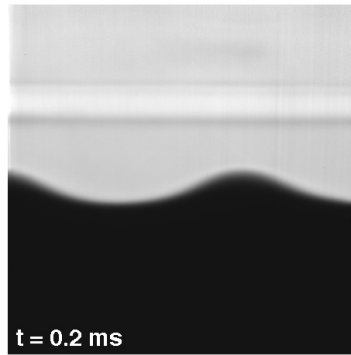
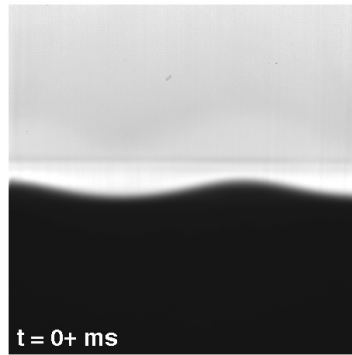
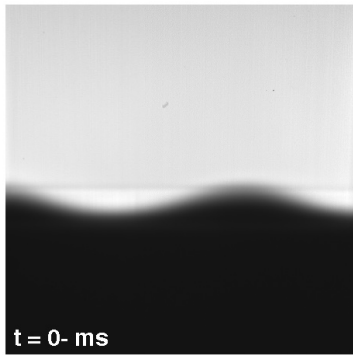
J.W. Jacobs and V.V. Krivets

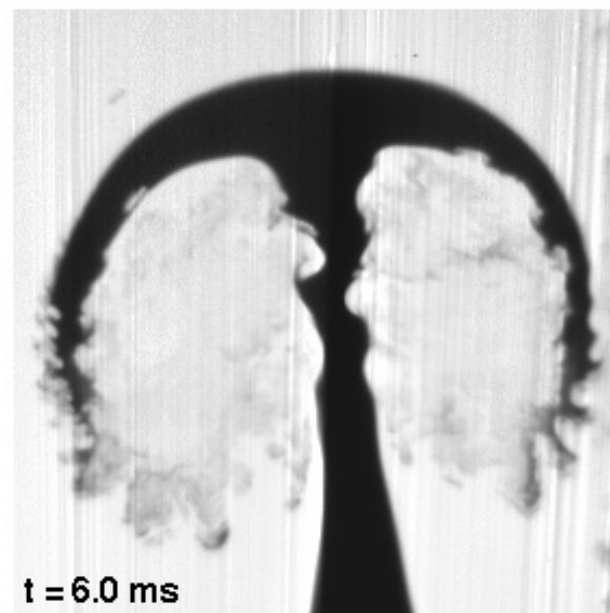
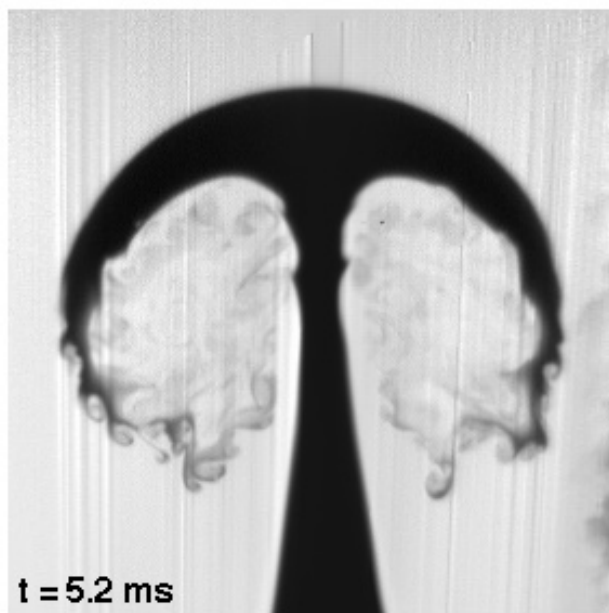
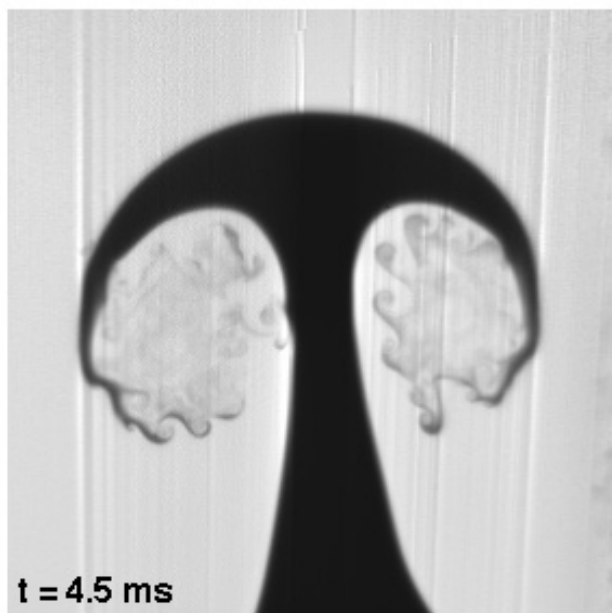
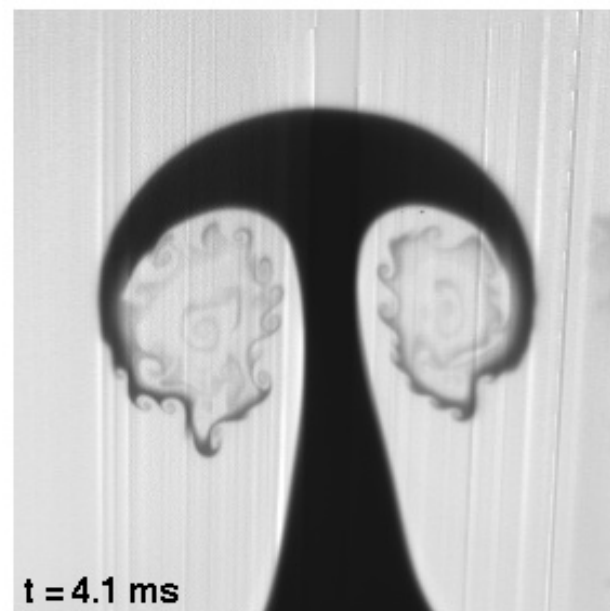
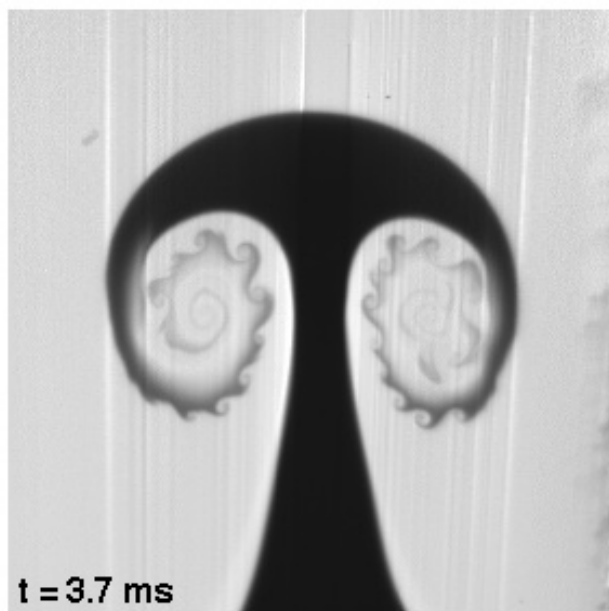
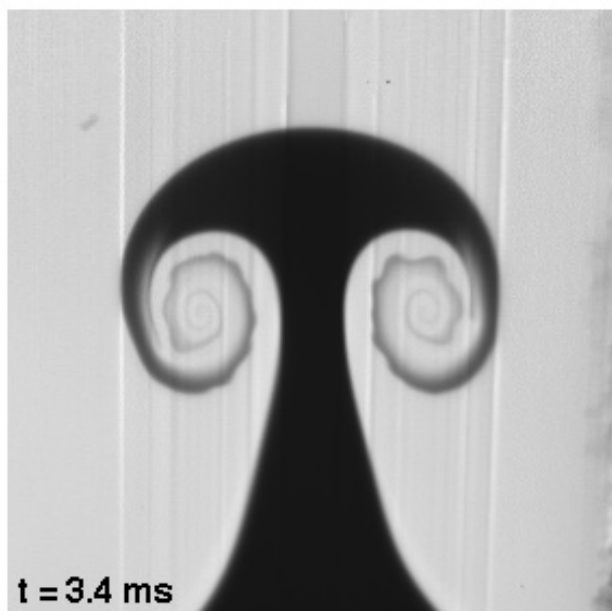
Department of Aerospace and Mechanical Engineering  
University of Arizona

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## Reynolds Number

$$\text{Re} = \frac{U\ell}{\nu}$$

Let,  $\ell \sim h = 2a$ ,  $U \sim \frac{dh}{dt}$ ,  $\nu$  = average viscosity

At time of transition,

$$\text{Re} \sim \frac{h\dot{h}}{\nu} = 50,000$$

Alternatively let,  $\text{Re} = \frac{U\ell}{\nu}$

Linear stability theory gives,  $\frac{U\ell}{\nu} = \frac{2}{\pi} \dot{a}_0$

Then,

$$\text{Re} = \frac{2}{\pi} \dot{a}_0 = 42,000$$

# Nonlinear Models

**Zhang & Sohn (1997)**

$$v_{b/s} = \frac{v_0}{1 + v_0 a_0 k^2 t + \max\left[0, a_0^2 k^2 - A^2 + \frac{1}{2}\right] v_0^2 k^2 t^2} + \frac{A v_0^2 k t}{1 + 2 k^2 a_0 v_0 t + 4 k^2 v_0^2 \left[ a_0^2 k^2 + \frac{1}{3} (1 - A^2) \right] t^2}$$

when  $A^2 = \frac{1}{2} + a_0^2 k^2$

$$v = \frac{1}{2} (v_b + v_s) \quad \frac{v_0}{\left[ a_0^2 k^2 - A^2 + \frac{1}{2} \right] v_0^2 k^2 t^2} \quad \frac{1}{t^2} \quad \text{as } t$$

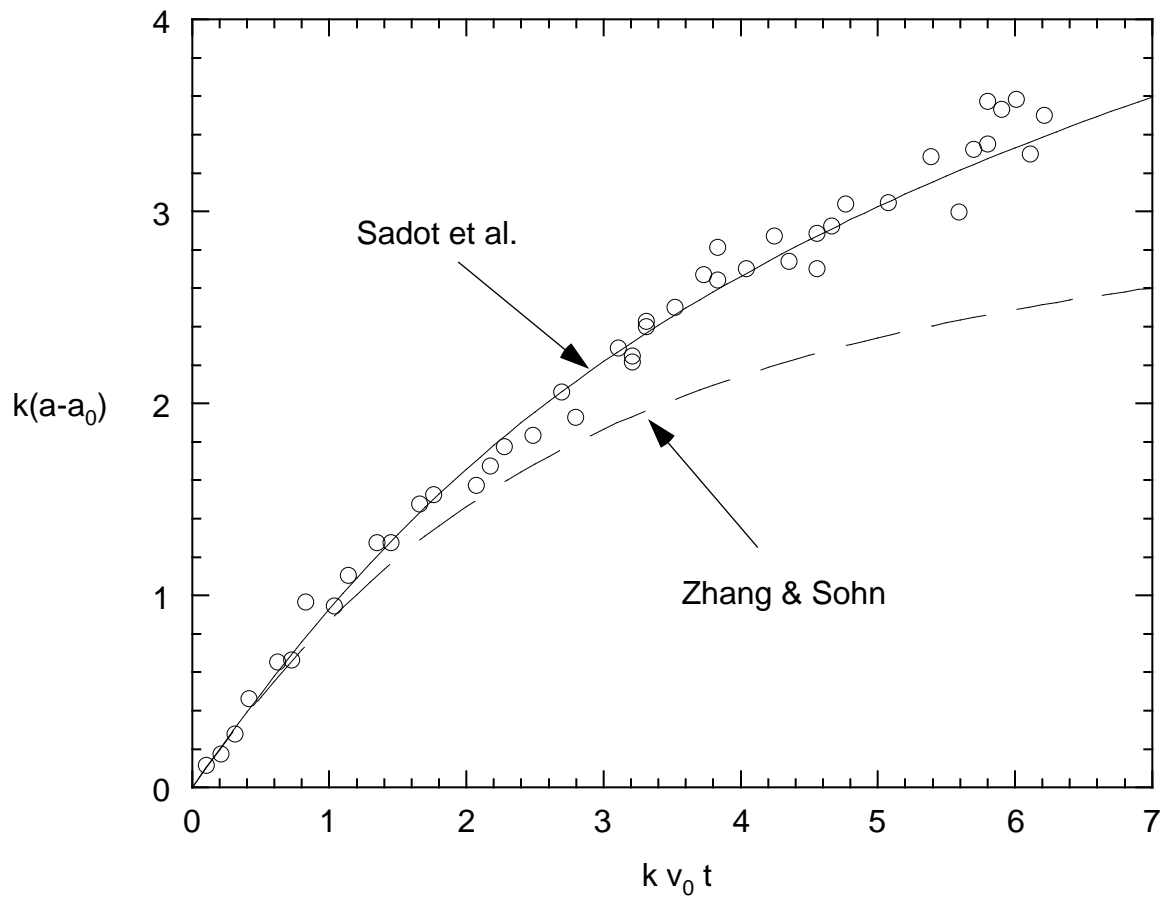
**Sadot el al. (1998)**

$$v_{b/s} = \frac{v_0 (1 + v_0 k t)}{1 + (1 \pm A) v_0 k t + \frac{1 \pm A}{1 + A} \frac{1}{2 C} v_0^2 k^2 t^2}$$

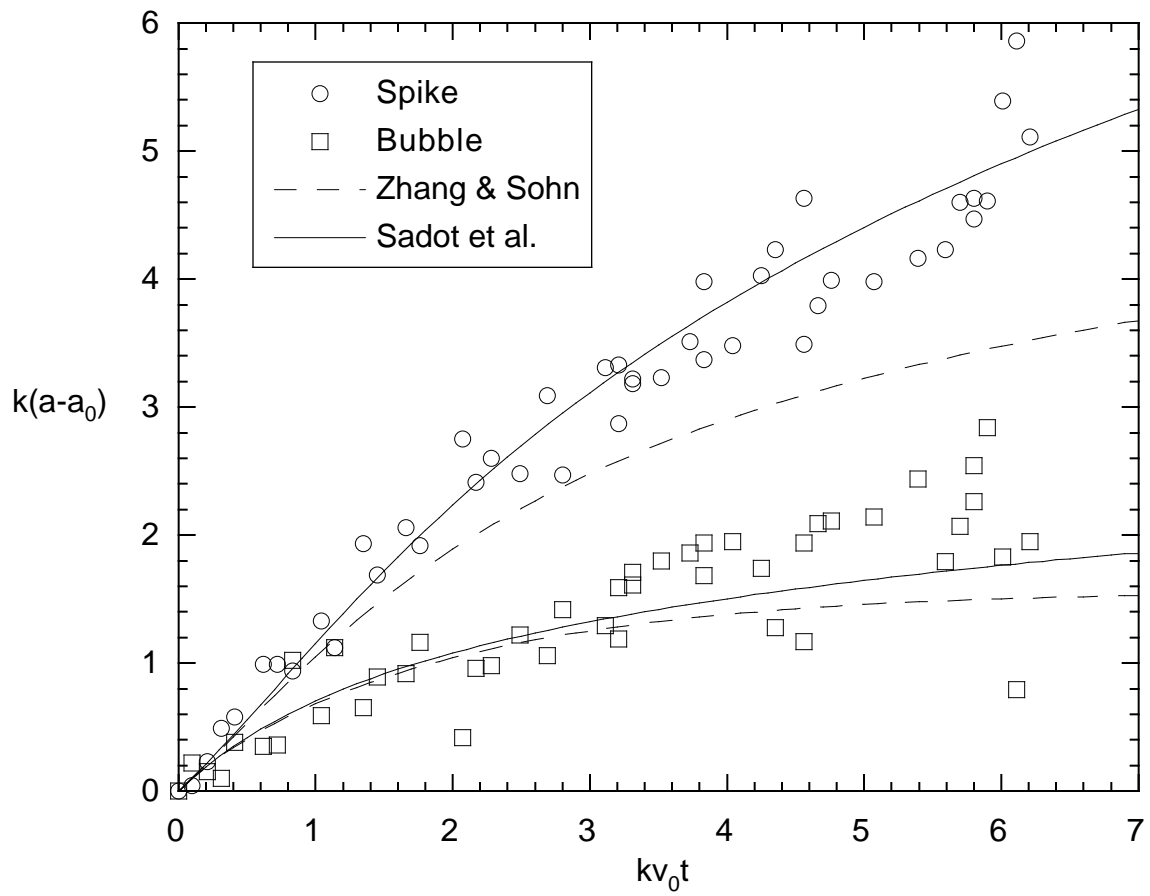
where  $C = \begin{matrix} 1/3 & A & 0.5 \\ 1/2 & A & 0 \end{matrix}$

$$v = \frac{1}{1 - A} \frac{C}{t} \frac{1}{t} \quad \text{as } t$$

### Late Time Growth

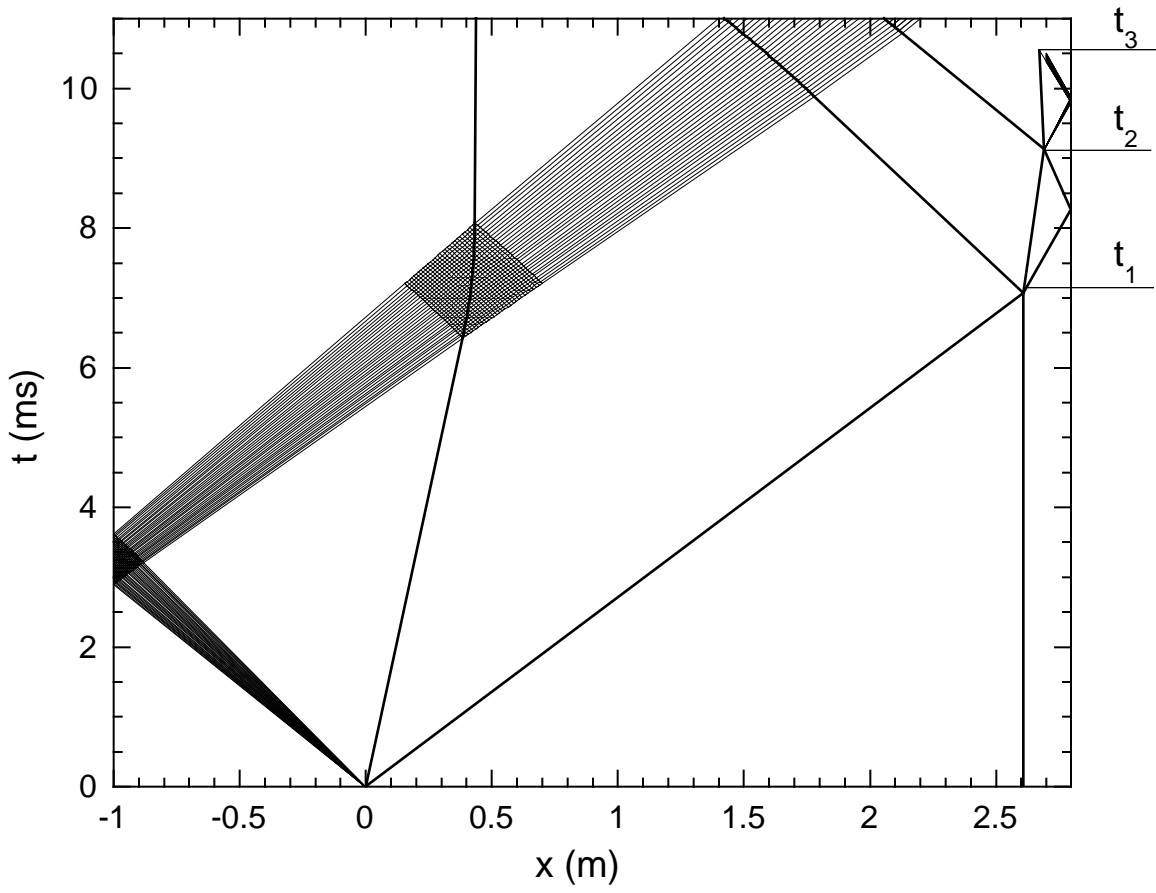


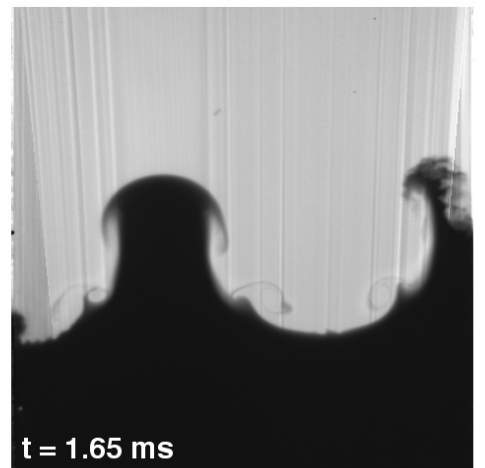
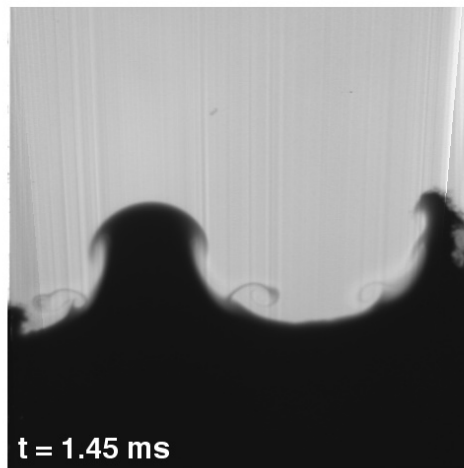
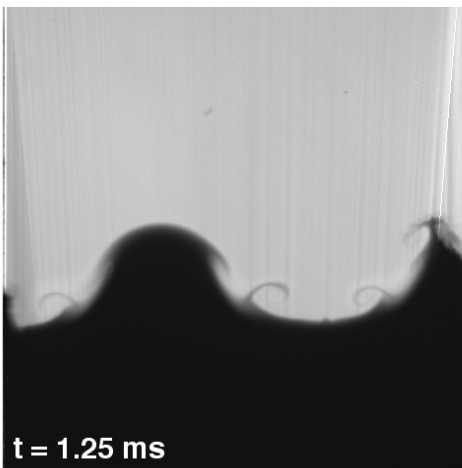
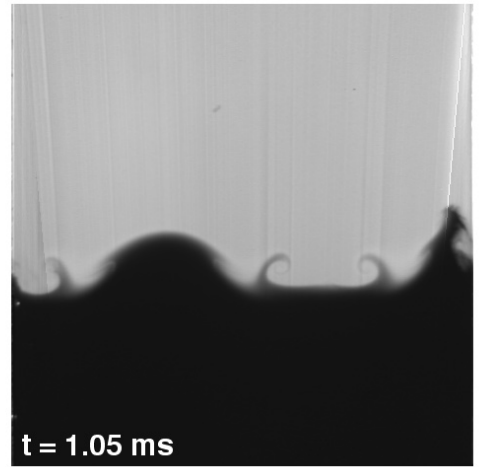
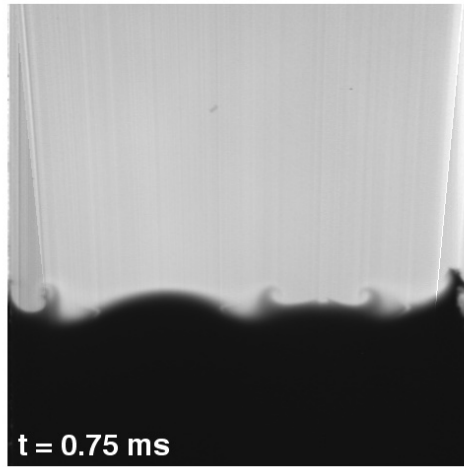
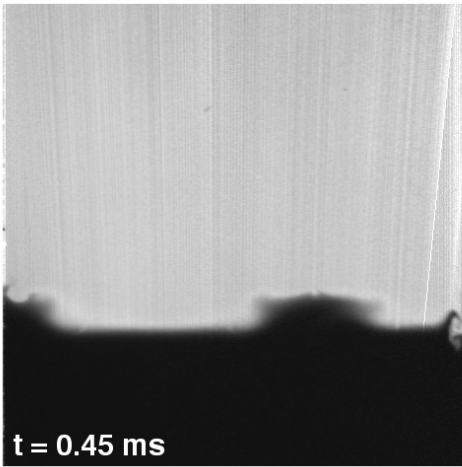
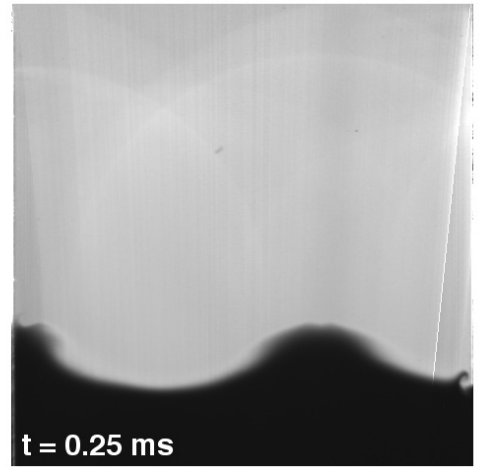
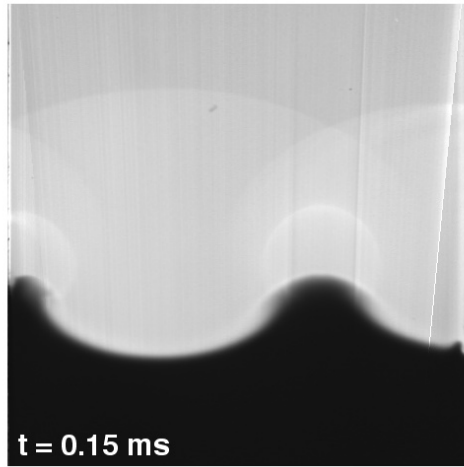
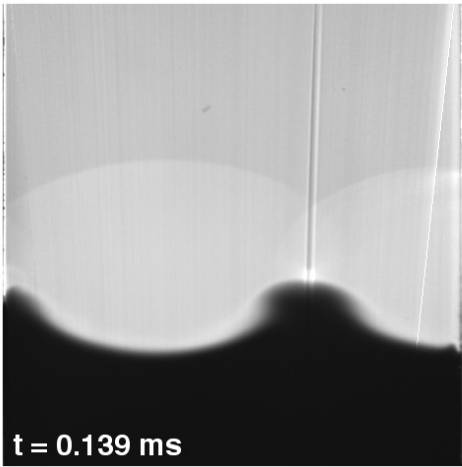
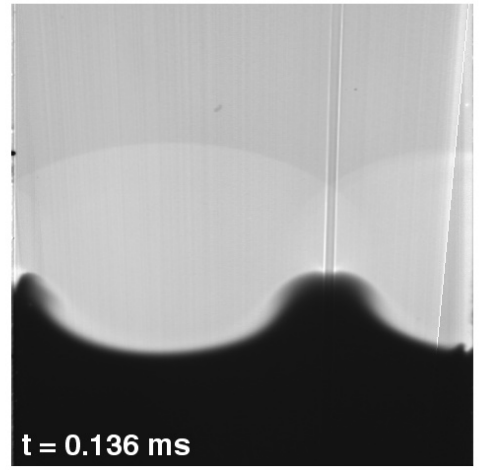
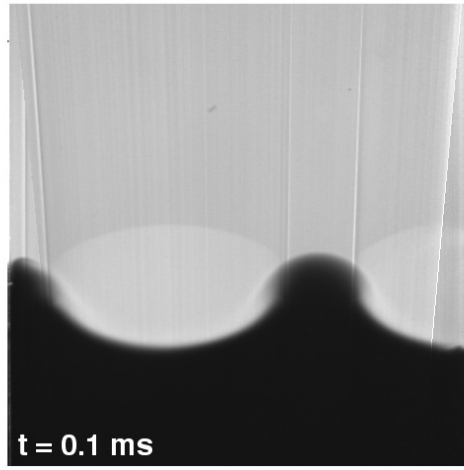
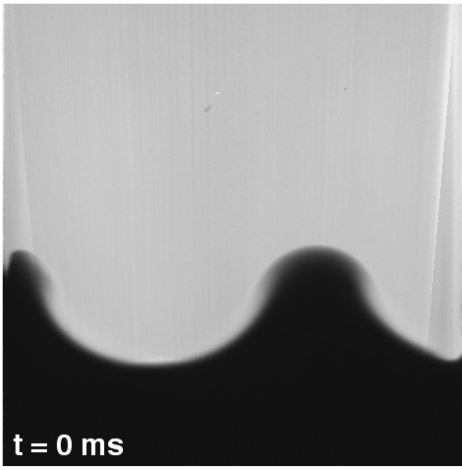
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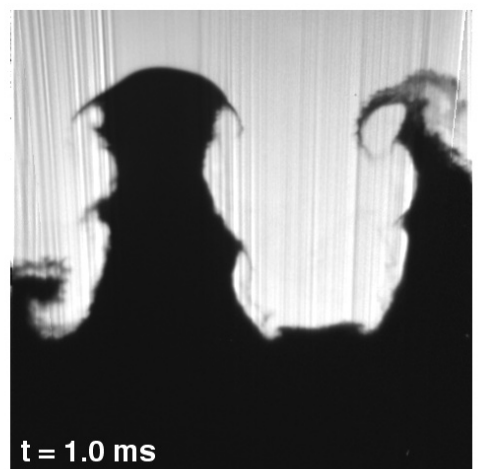
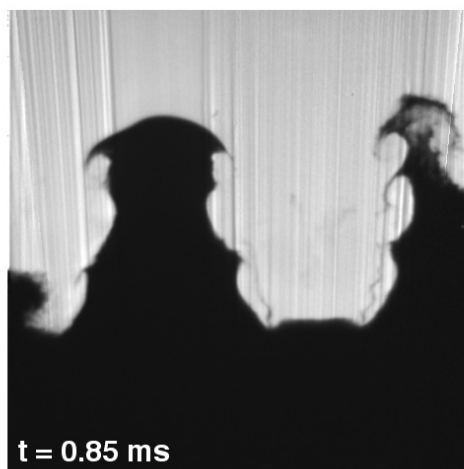
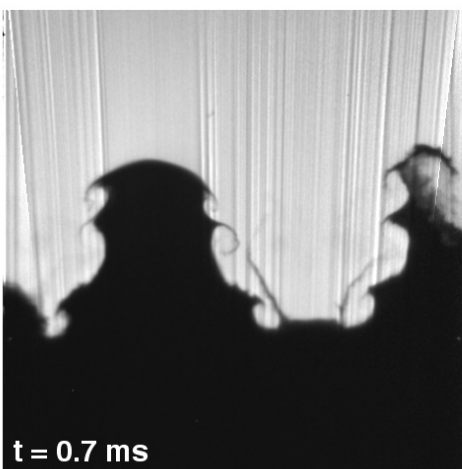
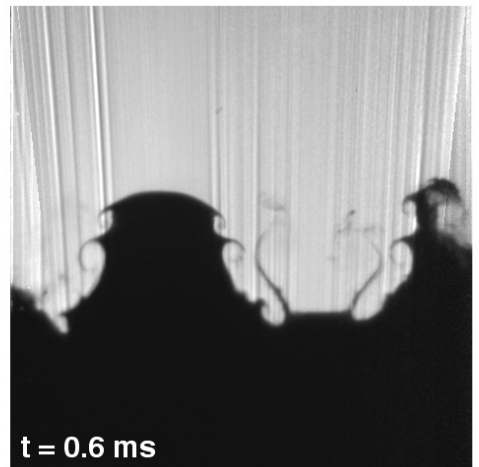
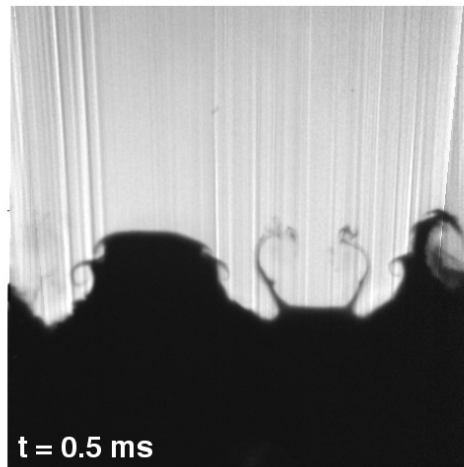
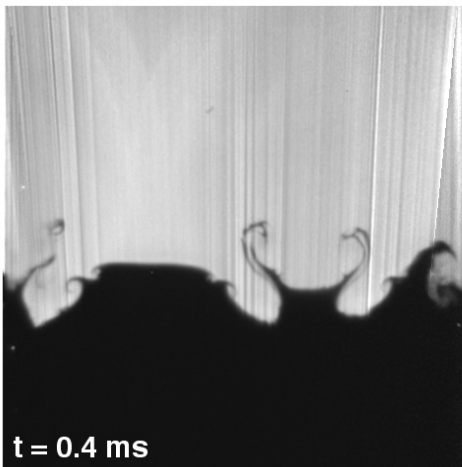
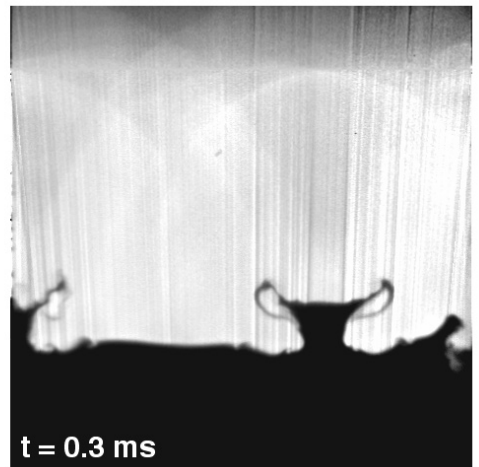
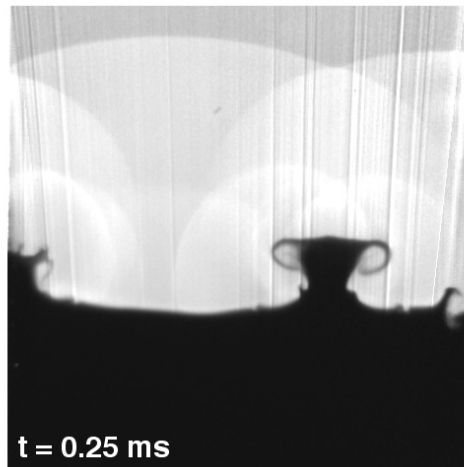
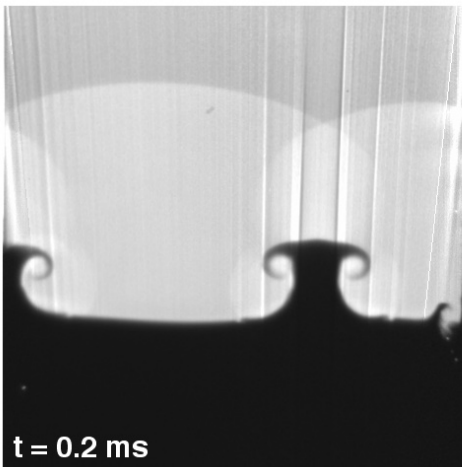
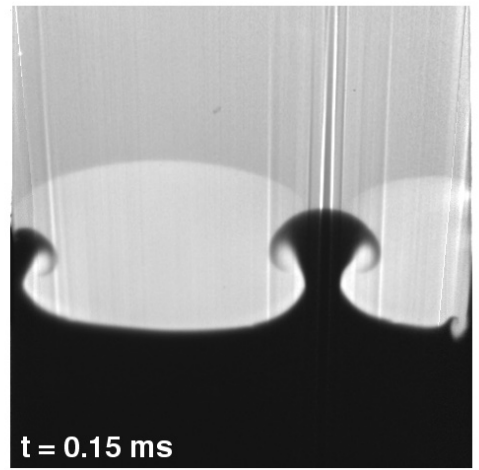
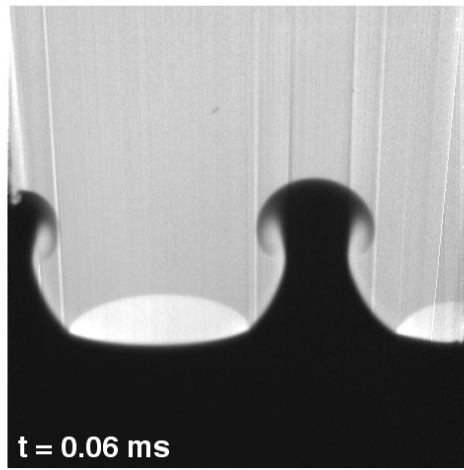
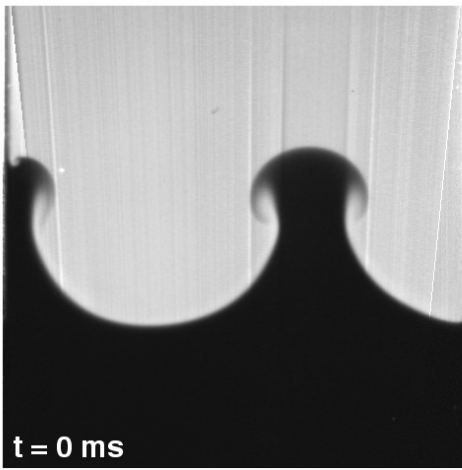


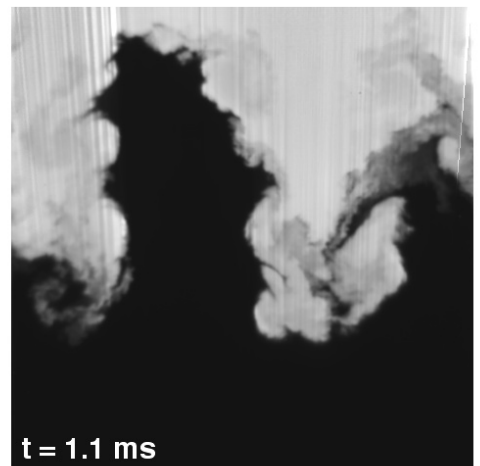
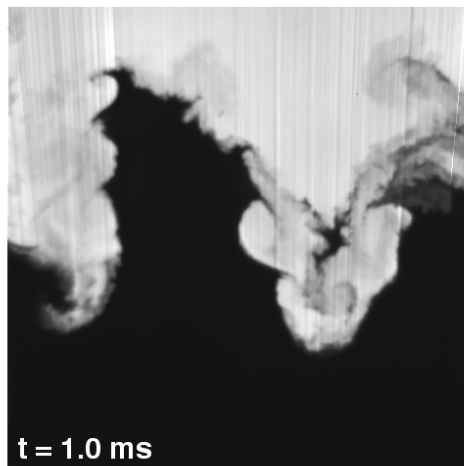
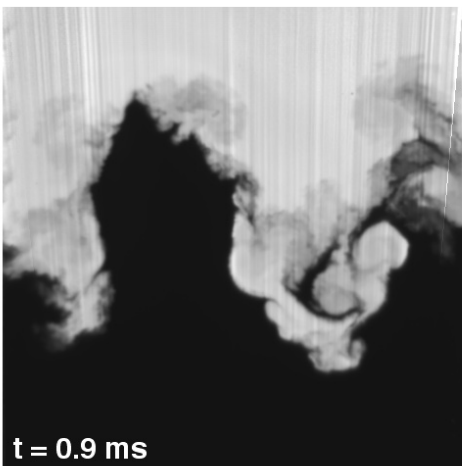
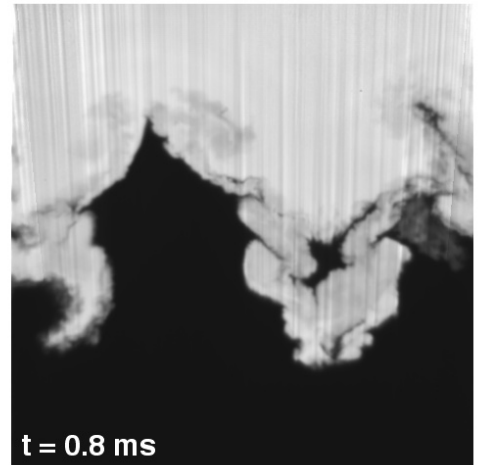
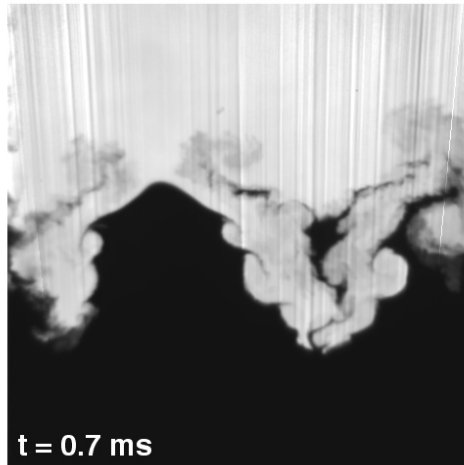
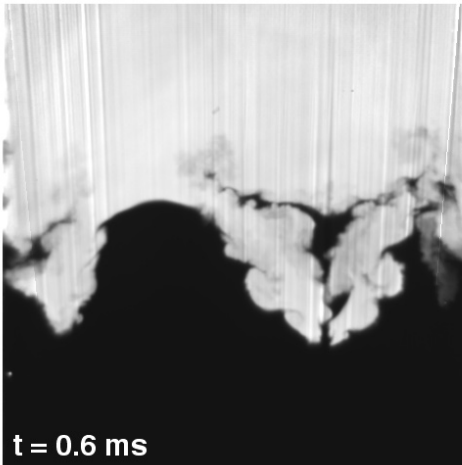
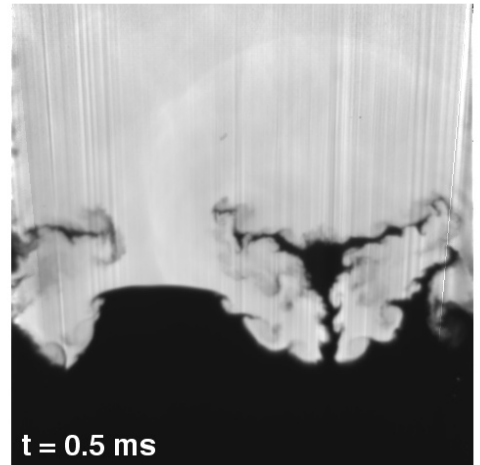
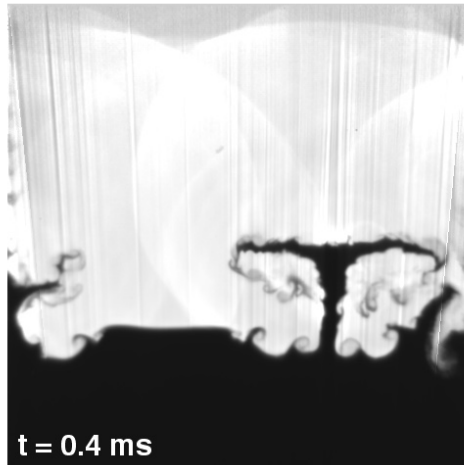
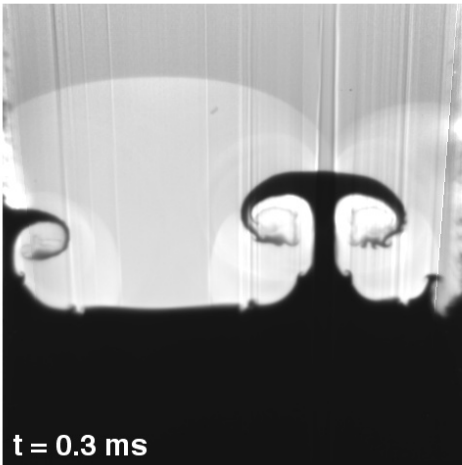
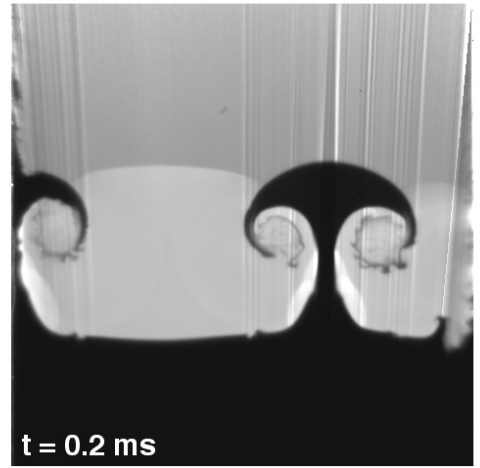
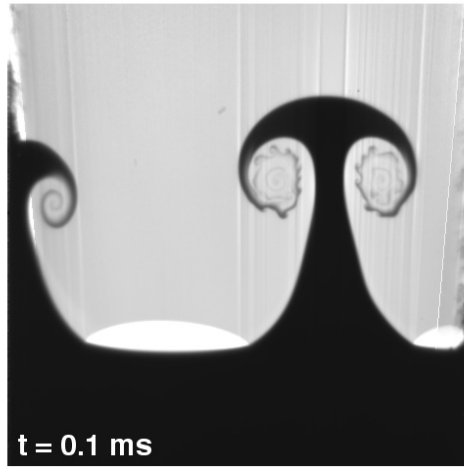
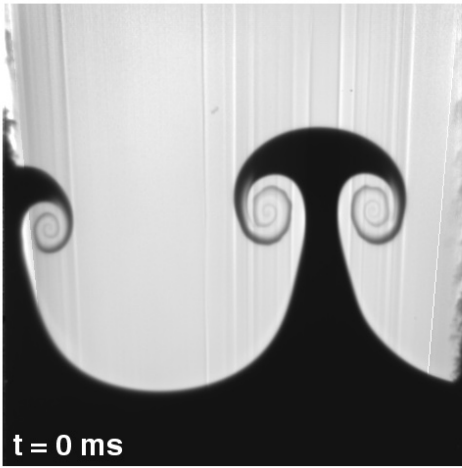


Wave Diagram;  $M = 1.11$ ;  $A = 0.656$

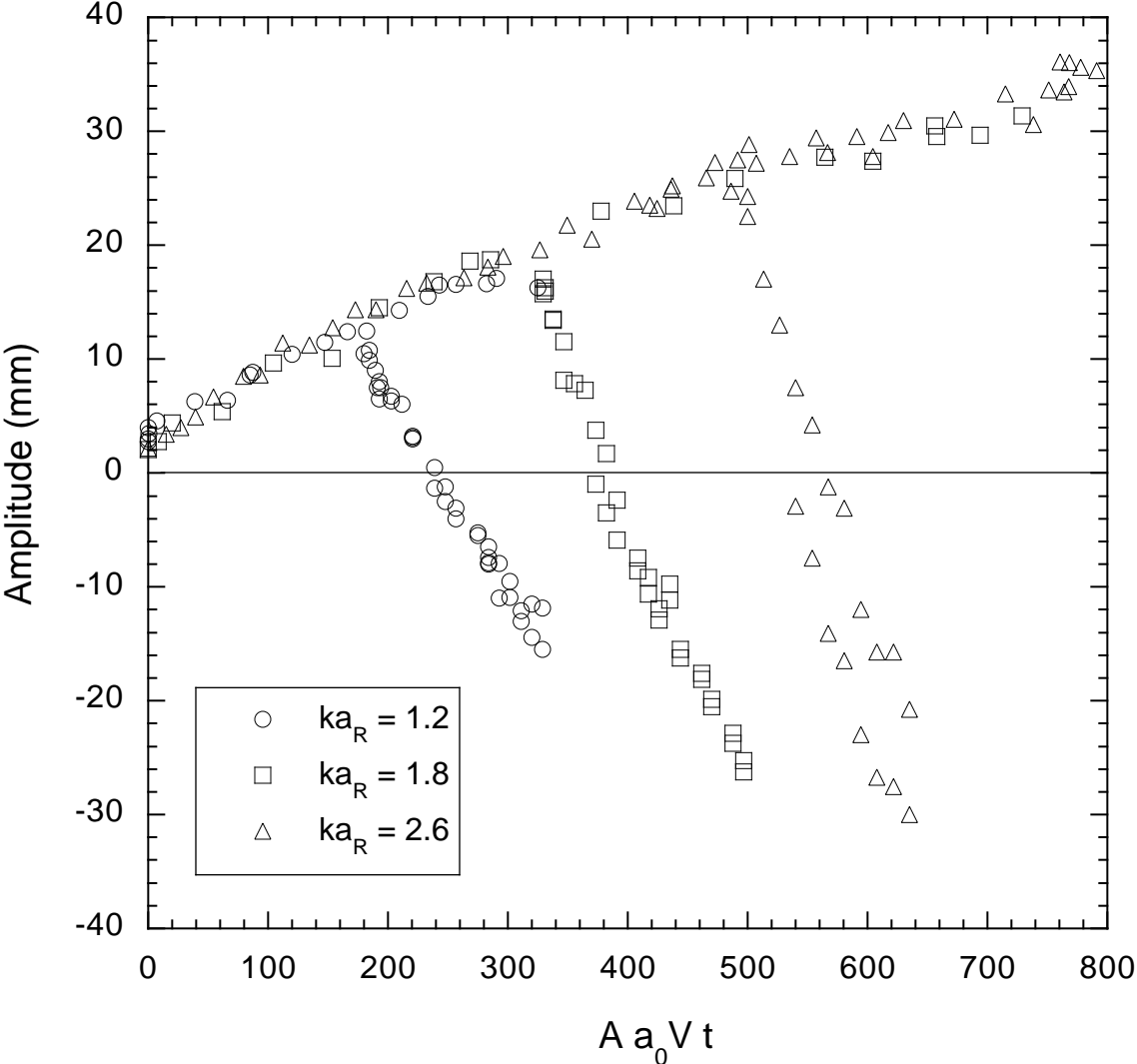




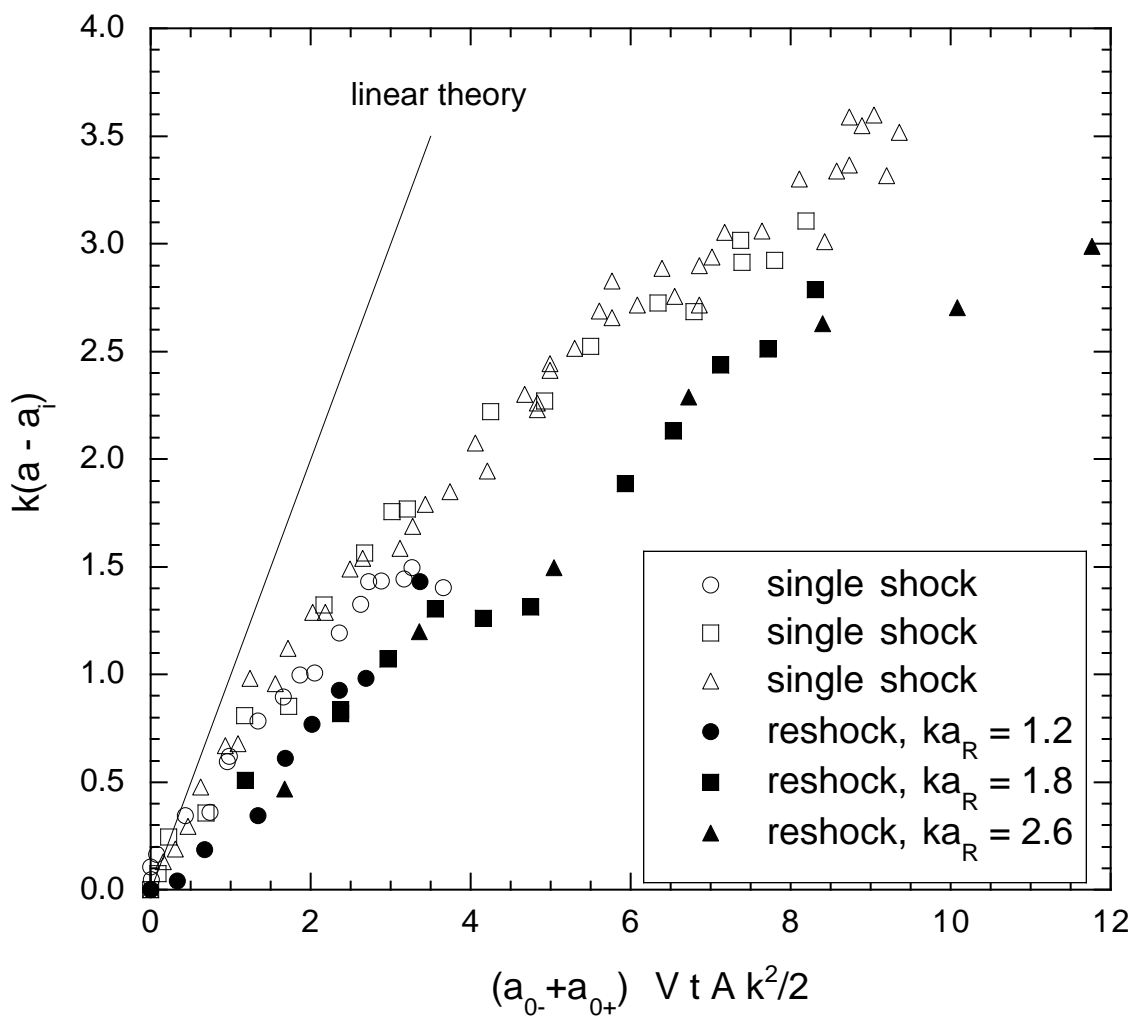




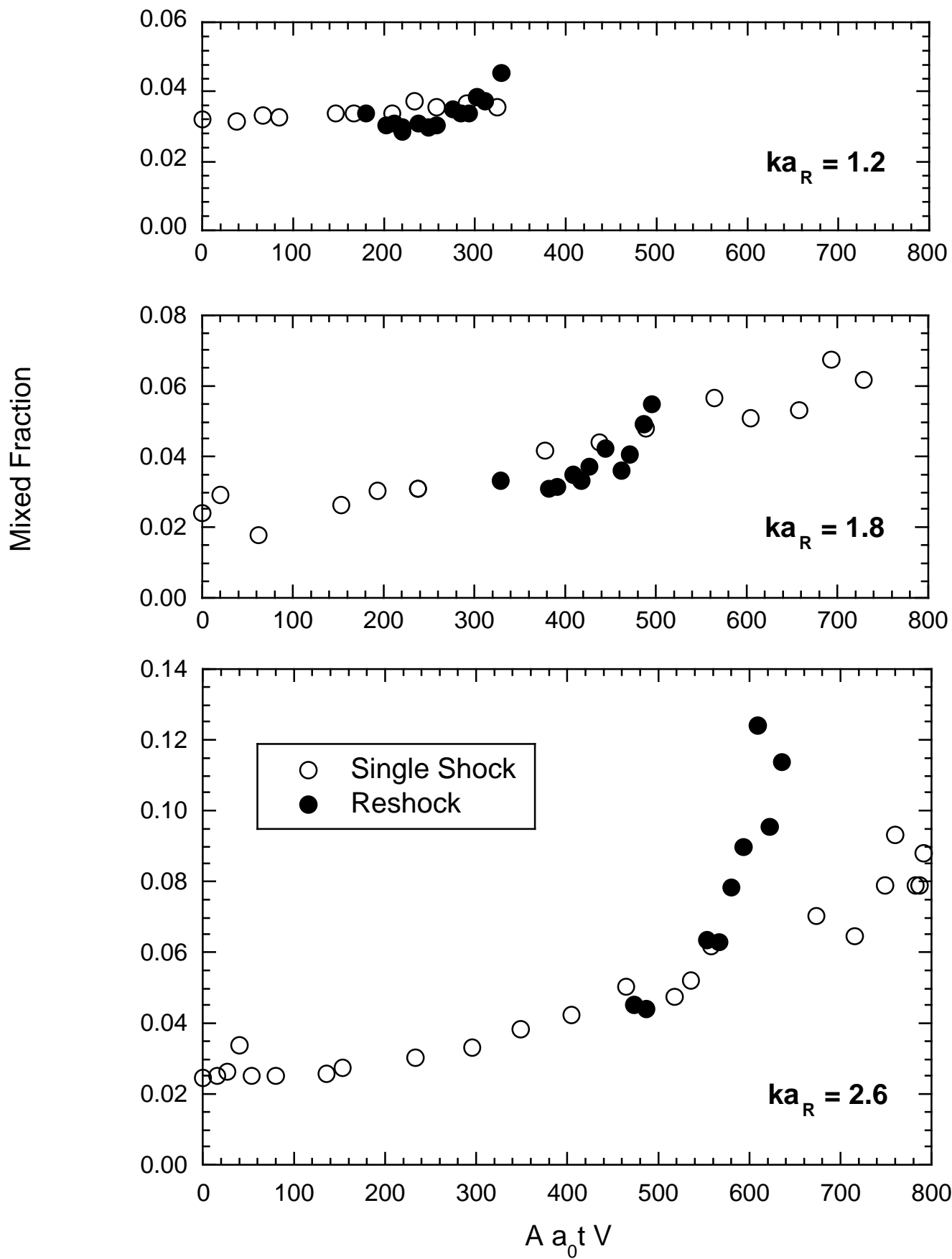
### Reshock Amplitude



### Growth Rate after Reshock



# 10% to 70% Mixed Fraction



## Conclusions

- ❑ Single-mode RM experiments show a transition to turbulence in the vortex cores at  $Re \approx 50,000$ .
- ❑ Late-time amplitude measurements show excellent agreement with the model of Sadot et al. (1998).
- ❑ Penetration depth measurements for the three reshock conditions collapse when plotted in dimensionless form and have a growth rate approximately  $1/3$  that given by linear stability theory.
- ❑ Reshock at early stages of the instability show similar evolution to that of single interaction with little increase in the mixing rate.
- ❑ Reshock at late stages of the instability produces a complex interface pattern and a significant increase in the mixing rate.